

Parton distribution functions and machine learning

P. Zurita

in collaboration with **D. Rentería-Estrada**, **R. Hernández-Pinto** and **G. Sborlini**



Universität Regensburg



Outlook

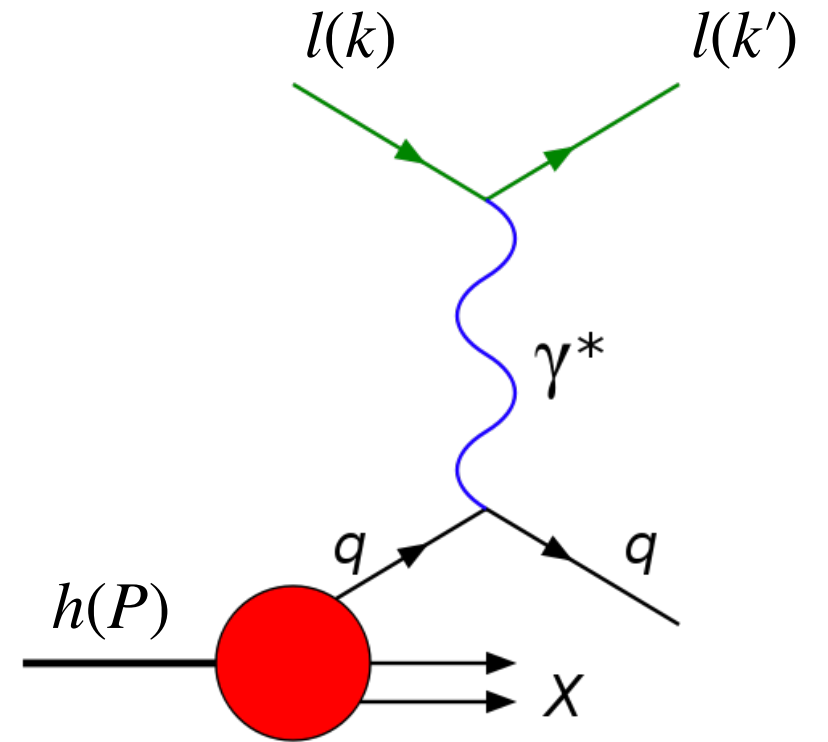
- ▶ Parton distributions functions:
 - ▶ Some history.
 - ▶ How to compute them.
 - ▶ Technical details.
- ▶ Using machine learning to learn about PDFs:
 - ▶ Motivation.
 - ▶ Accessing momentum fractions in p+p with ML (several methods and results).
- ▶ Going greener with ML (preliminary results).
- ▶ Summary

**“All” about parton
distribution functions**

1960s: Deep(l)y Inelastic Scattering

$$Q^2 = -(k - k')^2 \quad x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot k}$$

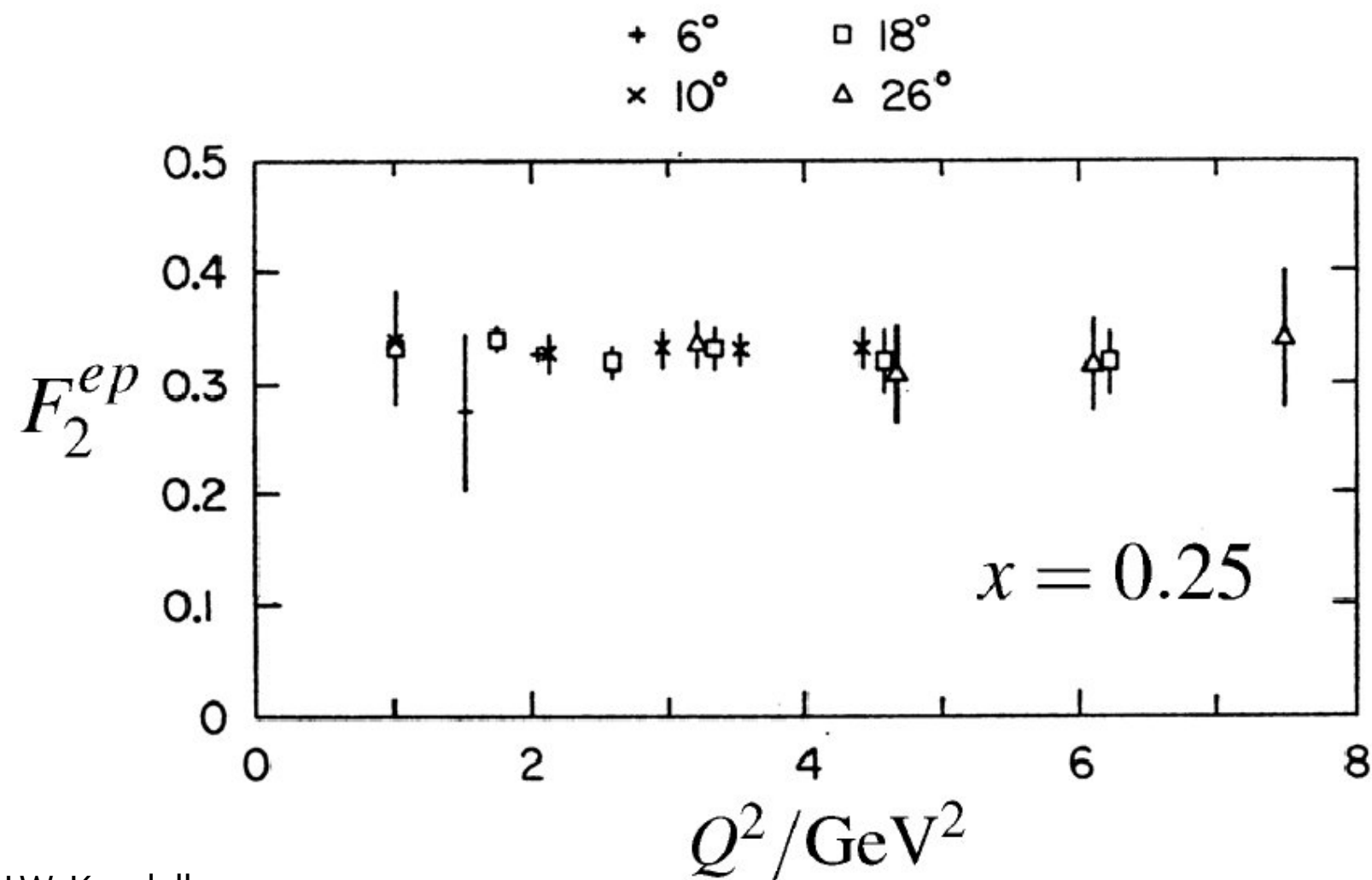
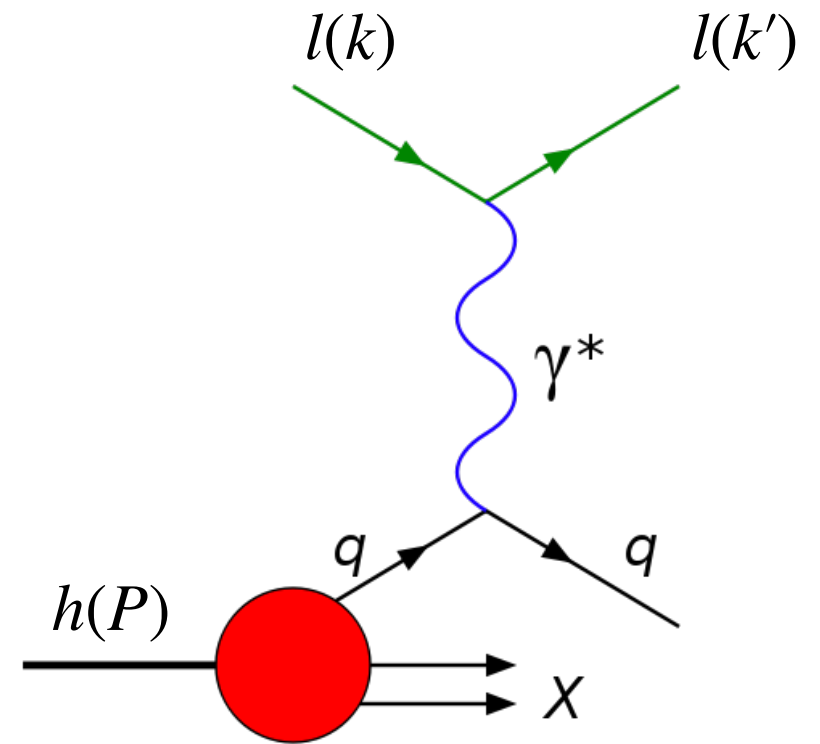
$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2)$$



1960s: Deep(l)y Inelastic Scattering

$$Q^2 = -(k - k')^2 \quad x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot k}$$

$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2)$$



- ▶ The scaling is expected if DIS is the incoherent scattering of partons (Feynman, 1969).

$$F_2^{LO}(x) = x \sum_{i=1}^{n_f} e_i^2 f_{i/h}(x)$$

- ▶ $f_{i/h}(x)$ is the **probability** density of finding the parton i inside the hadron h with x . These are called **Parton Distribution Functions** (PDFs).

- ▶ The scaling is expected if DIS is the incoherent scattering of partons (Feynman, 1969).

$$F_2^{LO}(x) = x \sum_{i=1}^{n_f} e_i^2 f_{i/h}(x)$$

- ▶ $f_{i/h}(x)$ is the **probability** density of finding the parton i inside the hadron h with x . These are called **Parton Distribution Functions** (PDFs).
- ▶ At higher orders in the expansion the scaling breaks down, a dependence on the renormalisation scale μ appears.

- ▶ The scaling is expected if DIS is the incoherent scattering of partons (Feynman, 1969).

$$F_2^{LO}(x) = x \sum_{i=1}^{n_f} e_i^2 f_{i/h}(x)$$

- ▶ $f_{i/h}(x)$ is the **probability** density of finding the parton i inside the hadron h with x . These are called **Parton Distribution Functions** (PDFs).
- ▶ At higher orders in the expansion the scaling breaks down, a dependence on the renormalisation scale μ appears.
- ▶ Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (**DGLAP**) equations give the evolution with the scale, while mixing the partons.

$$\mu^2 \frac{d}{d\mu^2} f_{a/h}(x, \mu) = \int_x^1 \frac{d\xi}{\xi} \sum_b P_{a/b}\left(\frac{x}{\xi}, \alpha_s(\mu)\right) f_{b/h}(\xi, \mu)$$

- ▶ These objects are related to the ones in the parton model and we still call them “PDFs”, but the probabilistic interpretation is not 100% accurate.

- ▶ These objects are related to the ones in the parton model and we still call them “PDFs”, but the probabilistic interpretation is not 100% accurate.
- ▶ The PDFs contain the long distance structure of hadrons.

- ▶ These objects are related to the ones in the parton model and we still call them “PDFs”, but the probabilistic interpretation is not 100% accurate.
- ▶ The PDFs contain the long distance structure of hadrons.
- ▶ They are **universal**, so any process that is an inclusive hard scattering can be written as

$$d\sigma^{DIS} = \sum_i d\sigma^{l+i \rightarrow l'} \otimes f_i$$

$$d\sigma^{DY} = \sum_{i,j} d\sigma^{i+j \rightarrow l+\bar{l}} \otimes f_i \otimes f_j$$

with the same PDFs.

- ▶ These objects are related to the ones in the parton model and we still call them “PDFs”, but the probabilistic interpretation is not 100% accurate.
- ▶ The PDFs contain the long distance structure of hadrons.
- ▶ They are **universal**, so any process that is an inclusive hard scattering can be written as

$$d\sigma^{DIS} = \sum_i d\sigma^{l+i \rightarrow l'} \otimes f_i$$

$$d\sigma^{DY} = \sum_{i,j} d\sigma^{i+j \rightarrow l+\bar{l}} \otimes f_i \otimes f_j$$

with the same PDFs.

Without PDFs there is no prediction!

- ▶ These objects are related to the ones in the parton model and we still call them “PDFs”, but the probabilistic interpretation is not 100% accurate.
- ▶ The PDFs contain the long distance structure of hadrons.
- ▶ They are **universal**, so any process that is an inclusive hard scattering can be written as

$$d\sigma^{DIS} = \sum_i d\sigma^{l+i \rightarrow l'} \otimes f_i$$

$$d\sigma^{DY} = \sum_{i,j} d\sigma^{i+j \rightarrow l+\bar{l}} \otimes f_i \otimes f_j$$

with the same PDFs.

Without PDFs there is no prediction!

I will adhere the KISS principle and stick to collinear PDFs.

- ▶ PDFs have formal definitions in terms of operators, for example

$$f_{j/h}(x, \mu) = \frac{1}{4\pi} \int dy^- e^{-ixP^+y^-} \langle P^+, \vec{0}_T | \bar{\psi}_j(0, y^-, \vec{0}_T) \gamma^+ \mathcal{O} \psi_j(0, 0, \vec{0}_T) | P^+, \vec{0}_T \rangle_{\bar{M}S}$$

$$\mathcal{O} = \mathcal{P} \exp \left(ig \int_0^{y^-} dz^- A_a^+(0, z^-, \vec{0}_T) t_a \right)$$

$$P^\pm = (P^0 \pm P^3)/\sqrt{2}$$

- ▶ PDFs have formal definitions in terms of operators, for example

$$f_{j/h}(x, \mu) = \frac{1}{4\pi} \int dy^- e^{-ixP^+y^-} \langle P^+, \vec{0}_T | \bar{\psi}_j(0, y^-, \vec{0}_T) \gamma^+ \mathcal{O} \psi_j(0, 0, \vec{0}_T) | P^+, \vec{0}_T \rangle_{\bar{MS}}$$

$$\mathcal{O} = \mathcal{P} \exp \left(ig \int_0^{y^-} dz^- A_a^+(0, z^-, \vec{0}_T) t_a \right)$$

$$P^\pm = (P^0 \pm P^3)/\sqrt{2}$$

- ▶ $|P\rangle$ is the state of a **hadron**.
- ▶ ψ is the field operator of a **quark**.

- ▶ PDFs have formal definitions in terms of operators, for example

$$f_{j/h}(x, \mu) = \frac{1}{4\pi} \int dy^- e^{-ixP^+y^-} \langle P^+, \vec{0}_T | \bar{\psi}_j(0, y^-, \vec{0}_T) \gamma^+ \mathcal{O} \psi_j(0, 0, \vec{0}_T) | P^+, \vec{0}_T \rangle_{\bar{M}S}$$

$$\mathcal{O} = \mathcal{P} \exp \left(ig \int_0^{y^-} dz^- A_a^+(0, z^-, \vec{0}_T) t_a \right)$$

$$P^\pm = (P^0 \pm P^3)/\sqrt{2}$$

- ▶ $|P\rangle$ is the state of a **hadron**.
- ▶ ψ is the field operator of a **quark**.
- ▶ So we can't compute these from first principles in pQCD, we must resort to phenomenology. We do **global** fits.

Steps for a “traditional” fit:

- ▶ Choose:
 - ▶ a factorisation scheme
 - ▶ an order in perturbation theory*
 - ▶ a starting scale Q_0 (so that above it pQCD is valid)
 - ▶ the data to be fitted
 - ▶ a heavy flavour scheme

Steps for a “traditional” fit:

- ▶ Choose:
 - ▶ a factorisation scheme
 - ▶ an order in perturbation theory*
 - ▶ a starting scale Q_0 (so that above it pQCD is valid)
 - ▶ the data to be fitted
 - ▶ a heavy flavour scheme
- ▶ Parametrise the quark and gluon distribution (or a combination of them):

$$xf_i(x, Q_0^2) = A_i x^{\alpha_i} (1 - x)^{\beta_i} P(x, c_i)$$

- ▶ Solve the DGLAP equations for the measured kinematics.
- ▶ Compute the hard cross sections for the observables.
- ▶ Convolute PDFs and partonic cross-sections.

- Compute this quantity and minimise it:

$$\sum_{i,j=1}^{N_{data}} \left[\sigma_{exp} - \sigma_{th} \right]_i C_{ij}^{-1} \left[\sigma_{exp} - \sigma_{th} \right]_j = \chi_{test}^2$$

- ▶ Compute this quantity and minimise it:

$$\sum_{i,j=1}^{N_{data}} \left[\sigma_{exp} - \sigma_{th} \right]_i C_{ij}^{-1} \left[\sigma_{exp} - \sigma_{th} \right]_j = \chi_{test}^2$$

- ▶ By the CLT, each term in the sum is distributed according to the square of a standard Gaussian.
- ▶ If we use d parameters, χ_{test}^2 follows a χ^2 distribution with $N_{data} - d$ degrees of freedom.
- ▶ $E[\chi_{test}^2] = N_{data} - d = N_{d.o.f.} \implies E[\chi_{test}^2]/N_{d.o.f.} = 1$

- ▶ Compute this quantity and minimise it:

$$\sum_{i,j=1}^{N_{data}} \left[\sigma_{exp} - \sigma_{th} \right]_i C_{ij}^{-1} \left[\sigma_{exp} - \sigma_{th} \right]_j = \chi_{test}^2$$

- ▶ By the CLT, each term in the sum is distributed according to the square of a standard Gaussian.
- ▶ If we use d parameters, χ_{test}^2 follows a χ^2 distribution with $N_{data} - d$ degrees of freedom.
- ▶ $E[\chi_{test}^2] = N_{data} - d = N_{d.o.f.} \implies E[\chi_{test}^2]/N_{d.o.f.} = 1$
- ▶ When we reach this point (the fit is “good enough”) we can be happy and store the final parameters.
- ▶ Use some method to estimate theoretical error bands.
- ▶ Create grids in x and Q^2 , and provide an interpolator for the grid. Make it publicly available (**LHAPDF**).

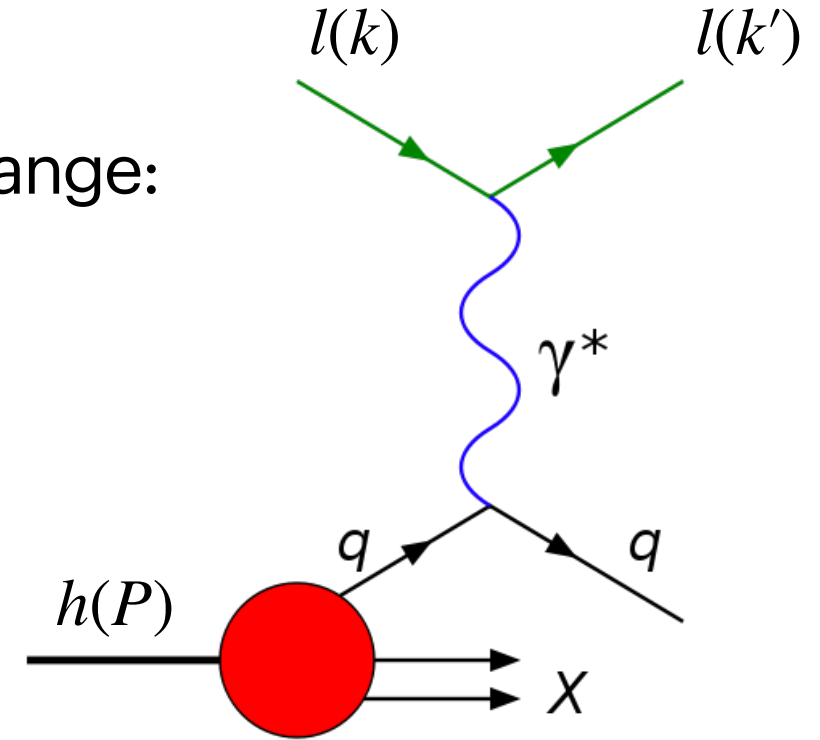
Some details

- ▶ The basic process is DIS with photon/Z boson exchange:

$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2)$$

$$F_2(x, Q^2) = \sum_{i=0}^{+\infty} \alpha_s^i(Q^2) \left[C_{2,q}^{i,NS} \otimes f_q^{NS} + C_{2,q}^{i,S} \otimes f_q^S + C_{2,g}^i \otimes f_g \right]$$

$$F_L(x, Q^2) = \sum_{i=1}^{+\infty} \alpha_s^i(Q^2) \left[C_{L,q}^{i,NS} \otimes f_q^{NS} + C_{L,q}^{i,S} \otimes f_q^S + C_{L,g}^i \otimes f_g \right]$$



Some details

- ▶ The basic process is DIS with photon/Z boson exchange:

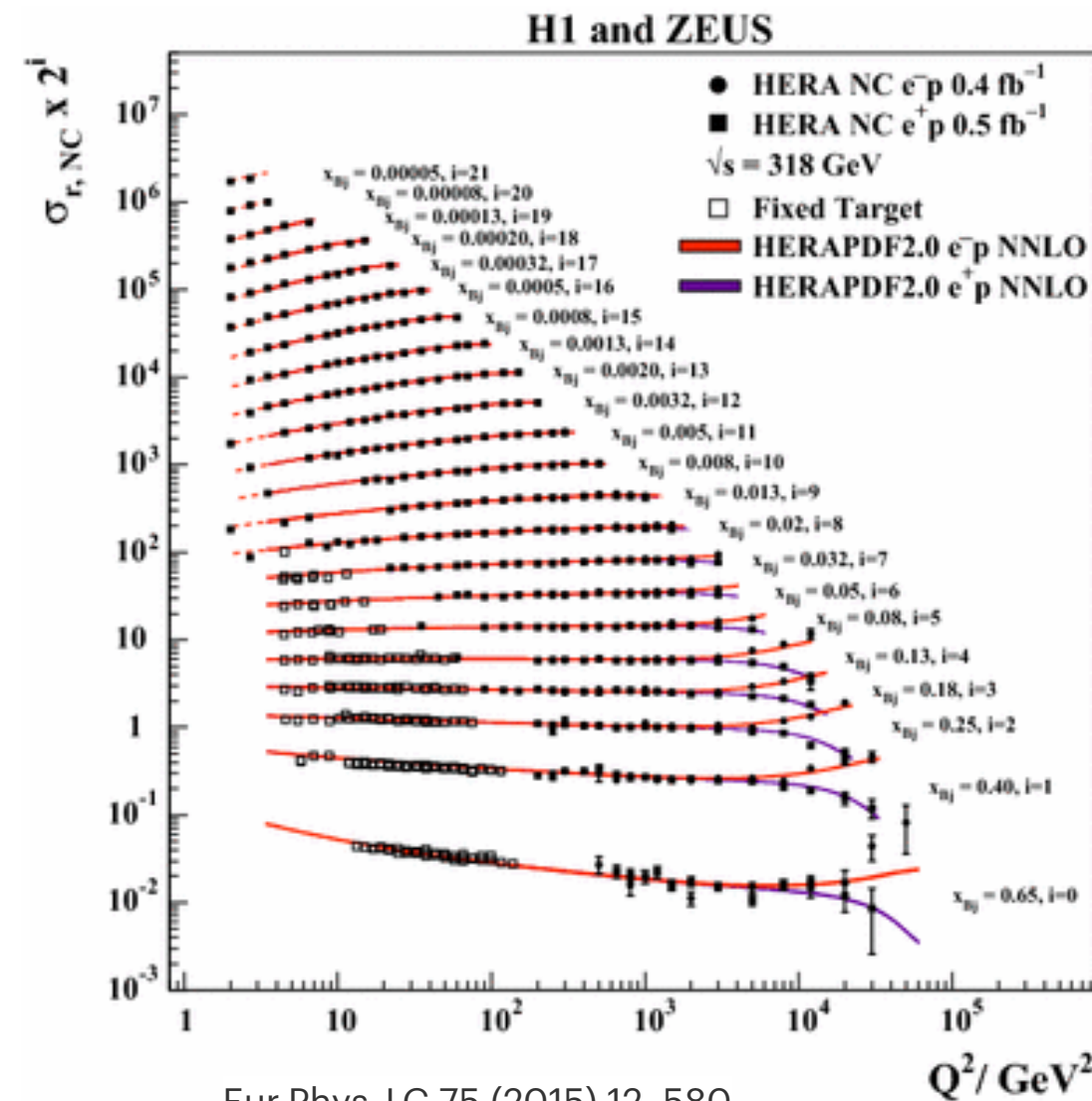
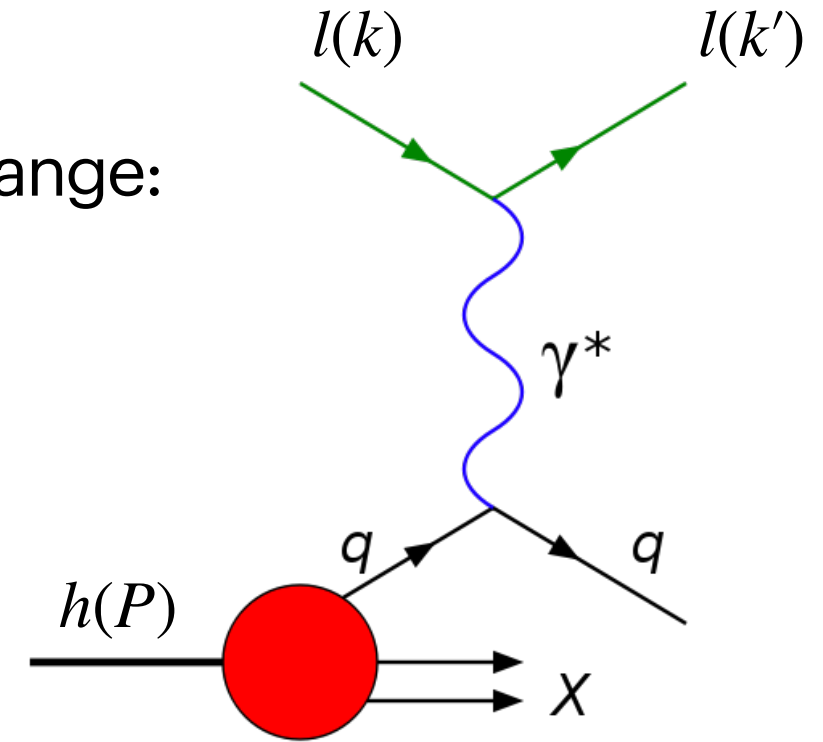
$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2)$$

$$F_2(x, Q^2) = \sum_{i=0}^{+\infty} \alpha_s^i(Q^2) \left[C_{2,q}^{i,NS} \otimes f_q^{NS} + C_{2,q}^{i,S} \otimes f_q^S + C_{2,g}^i \otimes f_g \right]$$

$$F_L(x, Q^2) = \sum_{i=1}^{+\infty} \alpha_s^i(Q^2) \left[C_{L,q}^{i,NS} \otimes f_q^{NS} + C_{L,q}^{i,S} \otimes f_q^S + C_{L,g}^i \otimes f_g \right]$$

$$f_q^{NS}(x, Q^2) = \sum_{i=1}^{n_f} e_i^2 f_i(x, Q^2)$$

$$f_q^S(x, Q^2) = \sum_{i=1}^{n_f} f_i(x, Q^2)$$

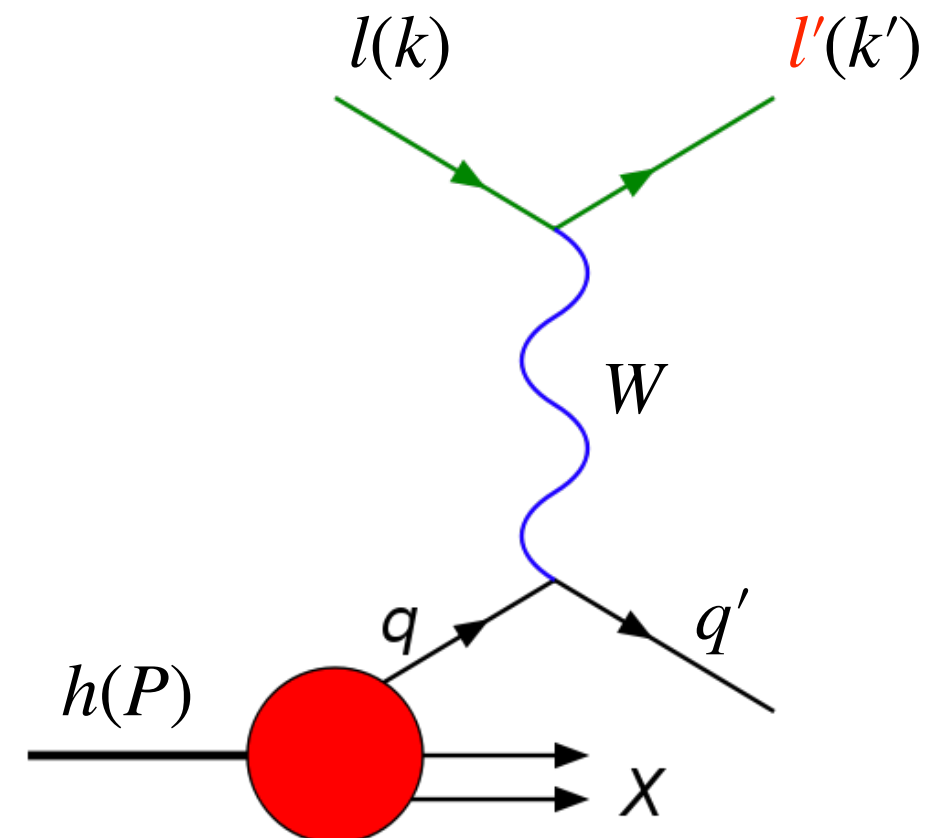


- ▶ For quarks, Neutral Current DIS is sensitive only to the non-singlet and singlet combination of the PDFs.

- ▶ For quarks, Neutral Current DIS is sensitive only to the non-singlet and singlet combination of the PDFs.
- ▶ We need to use other observables, such as Charged Current DIS to distinguish flavours.

$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2) \mp \frac{1 - (1 - y)^2}{1 + (1 - y)^2} xF_3$$

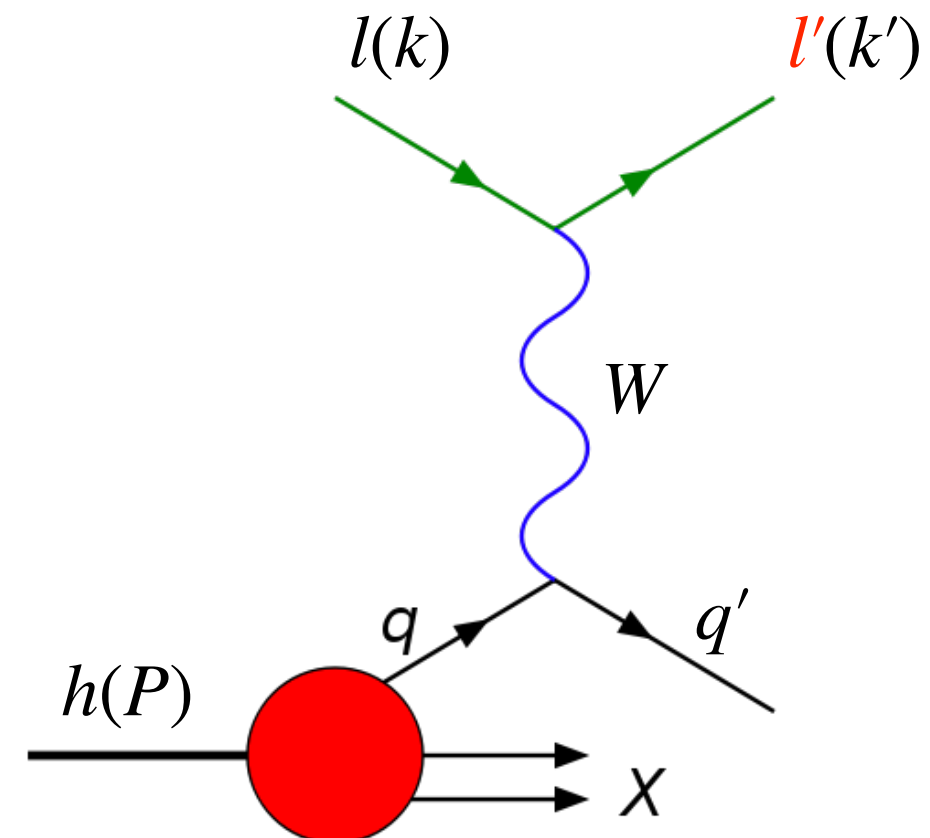
- ▶ xF_3 provides more PDF combinations, but it not enough (plus it is very complicated).



- ▶ For quarks, Neutral Current DIS is sensitive only to the non-singlet and singlet combination of the PDFs.
- ▶ We need to use other observables, such as Charged Current DIS to distinguish flavours.

$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2) \mp \frac{1 - (1 - y)^2}{1 + (1 - y)^2} xF_3$$

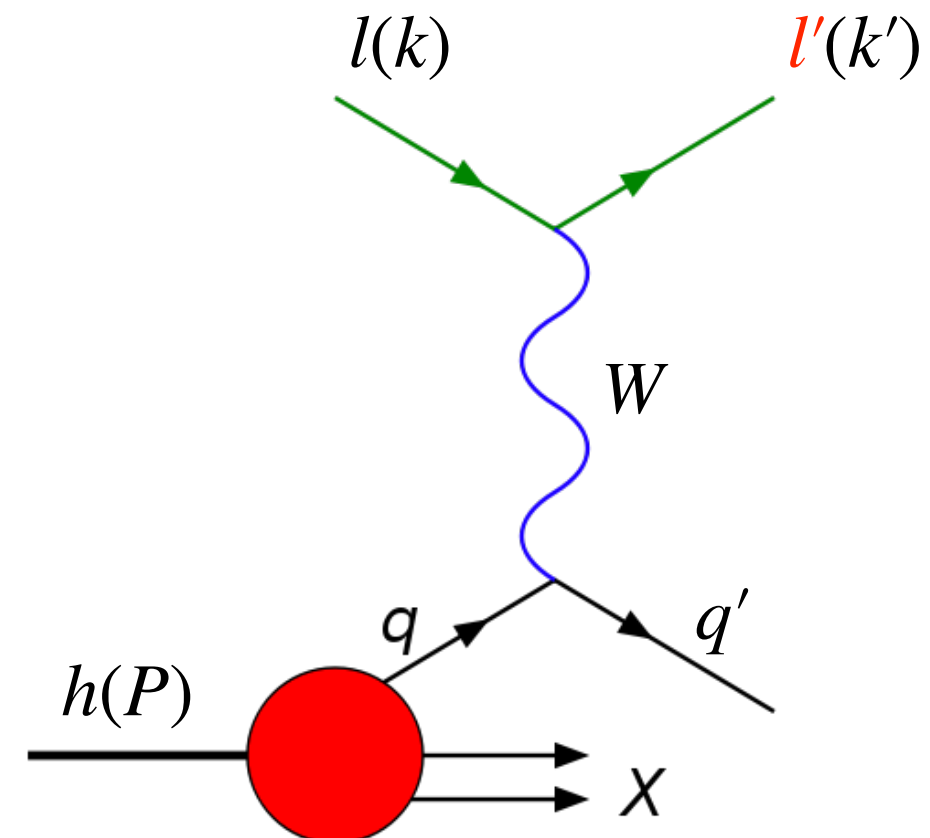
- ▶ xF_3 provides more PDF combinations, but it not enough (plus it is very complicated).
- ▶ We can use “neutrons”, but it is still to enough, due to the kinematic reach of the data...



- ▶ For quarks, Neutral Current DIS is sensitive only to the non-singlet and singlet combination of the PDFs.
- ▶ We need to use other observables, such as Charged Current DIS to distinguish flavours.

$$\frac{d^2\sigma}{dx dQ^2} \propto F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2) \mp \frac{1 - (1 - y)^2}{1 + (1 - y)^2} xF_3$$

- ▶ xF_3 provides more PDF combinations, but it not enough (plus it is very complicated).
- ▶ We can use “neutrons”, but it is still to enough, due to the kinematic reach of the data...
- ▶ We use **p+p**!



► Let us see for example Drell-Yan: $\frac{d^2\sigma}{dx_F dQ^2} = \frac{d^2\sigma^A}{dx_F dQ^2} + \frac{d^2\sigma^C}{dx_F dQ^2}$

$$\frac{d^2\sigma^A}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \frac{d^2\hat{\sigma}^A}{dx_F dQ^2} [f_i(t_1, Q^2)\bar{f}_i(t_2, Q^2) + \bar{f}_i(t_1, Q^2)f_i(t_2, Q^2)]$$

$$\frac{d^2\sigma^C}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \left[\frac{d^2\hat{\sigma}^C}{dx_F dQ^2} f_g(t_1, Q^2) [f_i(t_2, Q^2) + \bar{f}_i(t_2, Q^2)] + (1 \longleftrightarrow 2) \right]$$

► Let us see for example Drell-Yan: $\frac{d^2\sigma}{dx_F dQ^2} = \frac{d^2\sigma^A}{dx_F dQ^2} + \frac{d^2\sigma^C}{dx_F dQ^2}$

$$\frac{d^2\sigma^A}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \frac{d^2\hat{\sigma}^A}{dx_F dQ^2} [f_i(t_1, Q^2)\bar{f}_i(t_2, Q^2) + \bar{f}_i(t_1, Q^2)f_i(t_2, Q^2)]$$

$$\frac{d^2\sigma^C}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \left[\frac{d^2\hat{\sigma}^C}{dx_F dQ^2} f_g(t_1, Q^2) [f_i(t_2, Q^2) + \bar{f}_i(t_2, Q^2)] + (1 \longleftrightarrow 2) \right]$$

- Now it is quadratic in the PDFs, therefore harder to extract.

► Let us see for example Drell-Yan: $\frac{d^2\sigma}{dx_F dQ^2} = \frac{d^2\sigma^A}{dx_F dQ^2} + \frac{d^2\sigma^C}{dx_F dQ^2}$

$$\frac{d^2\sigma^A}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \frac{d^2\hat{\sigma}^A}{dx_F dQ^2} [f_i(t_1, Q^2)\bar{f}_i(t_2, Q^2) + \bar{f}_i(t_1, Q^2)f_i(t_2, Q^2)]$$

$$\frac{d^2\sigma^C}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \left[\frac{d^2\hat{\sigma}^C}{dx_F dQ^2} f_g(t_1, Q^2) [f_i(t_2, Q^2) + \bar{f}_i(t_2, Q^2)] + (1 \longleftrightarrow 2) \right]$$

- Now it is quadratic in the PDFs, therefore harder to extract.
- More importantly, beyond LO the x_1, x_2 **are not the momentum fractions** of the partons in the hard interaction!

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$

► Let us see for example Drell-Yan: $\frac{d^2\sigma}{dx_F dQ^2} = \frac{d^2\sigma^A}{dx_F dQ^2} + \frac{d^2\sigma^C}{dx_F dQ^2}$

$$\frac{d^2\sigma^A}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \frac{d^2\hat{\sigma}^A}{dx_F dQ^2} [f_i(t_1, Q^2)\bar{f}_i(t_2, Q^2) + \bar{f}_i(t_1, Q^2)f_i(t_2, Q^2)]$$

$$\frac{d^2\sigma^C}{dx_F dQ^2} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 \int_{x_1}^1 dt_1 \int_{x_2}^1 dt_2 \left[\frac{d^2\hat{\sigma}^C}{dx_F dQ^2} f_g(t_1, Q^2) [f_i(t_2, Q^2) + \bar{f}_i(t_2, Q^2)] + (1 \longleftrightarrow 2) \right]$$

- Now it is quadratic in the PDFs, therefore harder to extract.
- More importantly, beyond LO the x_1, x_2 **are not the momentum fractions** of the partons in the hard interaction!

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$

definition used beyond LO (🤔 but it is not true!)

Machine learning for PDFs

Motivation

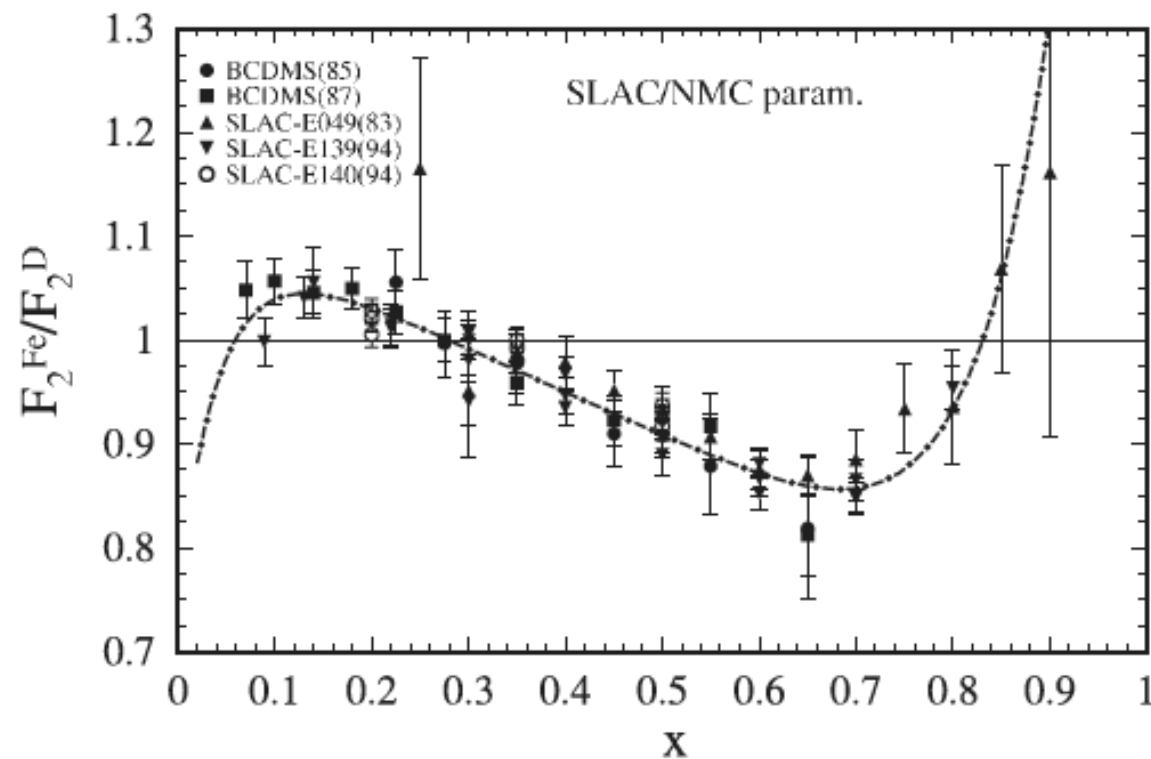
- ▶ Not being able to access the real x_1, x_2 from measuring the kinematics is not exclusive of the Drell-Yan process. It happens for all p+p collisions.
- ▶ And can have significant impact: e.g. in p+Pb collisions

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$

Motivation

- ▶ Not being able to access the real x_1, x_2 from measuring the kinematics is not exclusive of the Drell-Yan process. It happens for all p+p collisions.
- ▶ And can have significant impact: e.g. in p+Pb collisions

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$



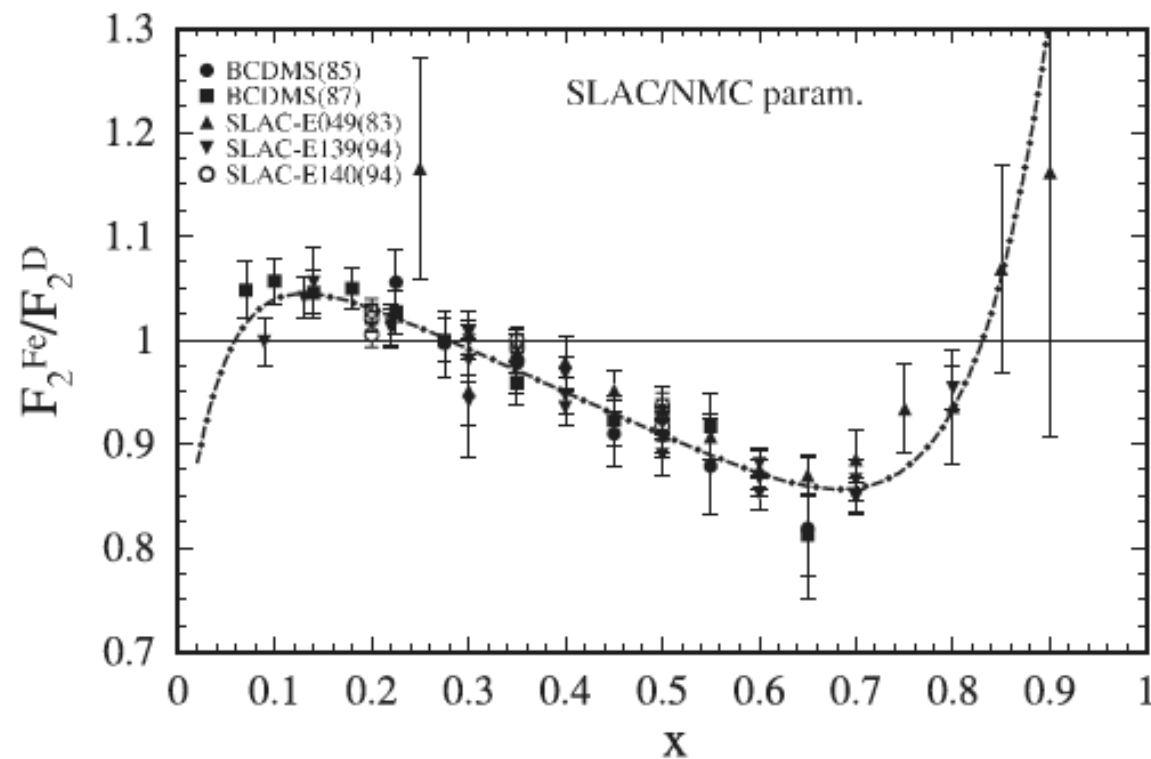
data from e+A DIS

Motivation

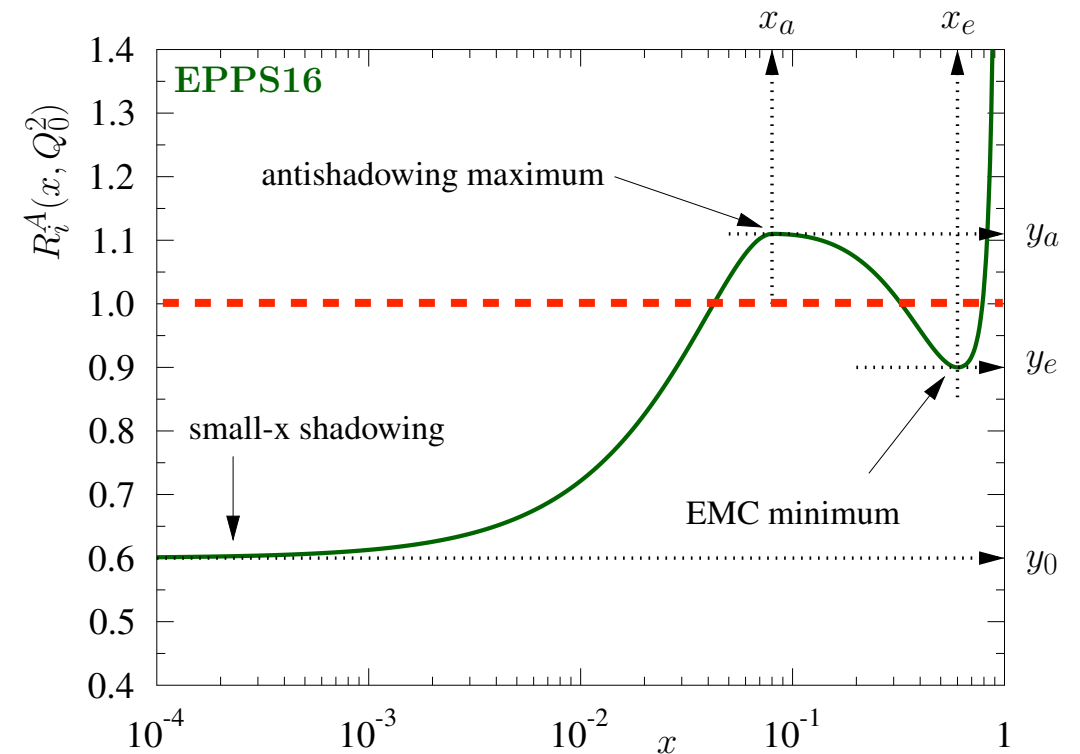
- ▶ Not being able to access the real x_1, x_2 from measuring the kinematics is not exclusive of the Drell-Yan process. It happens for all p+p collisions.

- ▶ And can have significant impact: e.g. in p+Pb collisions

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$



data from e+A DIS



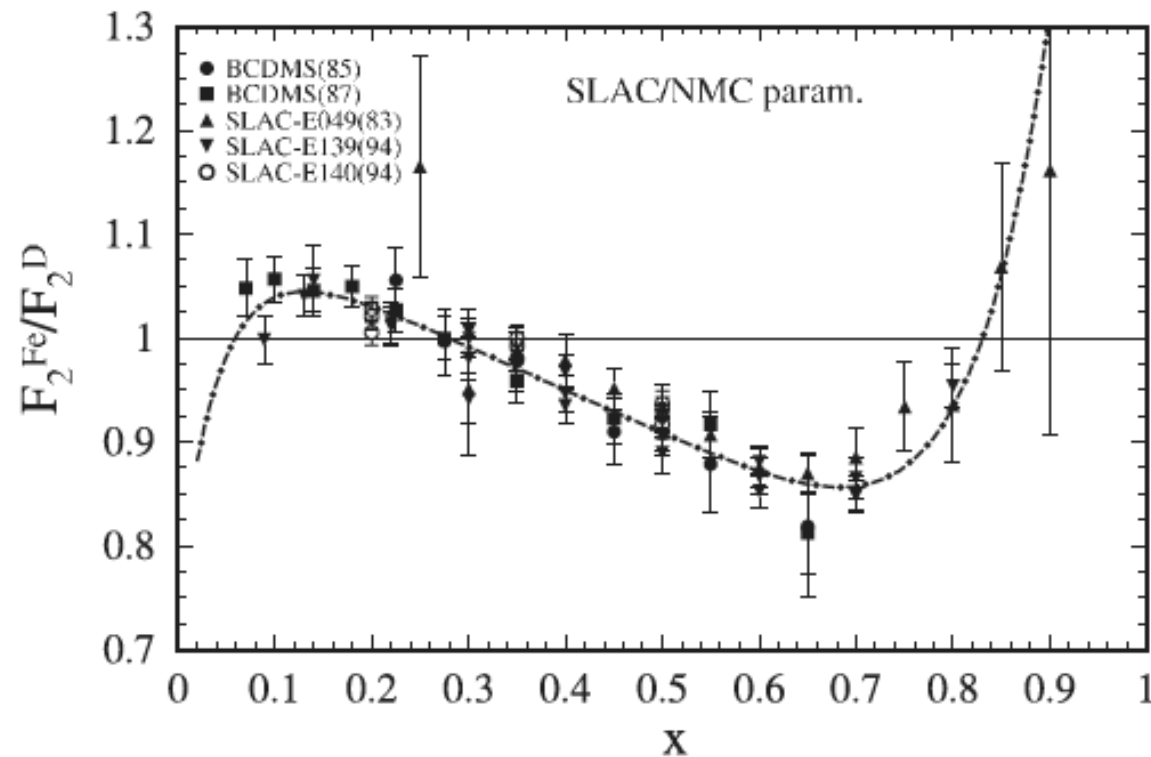
ratio of proton to nuclear PDF

Motivation

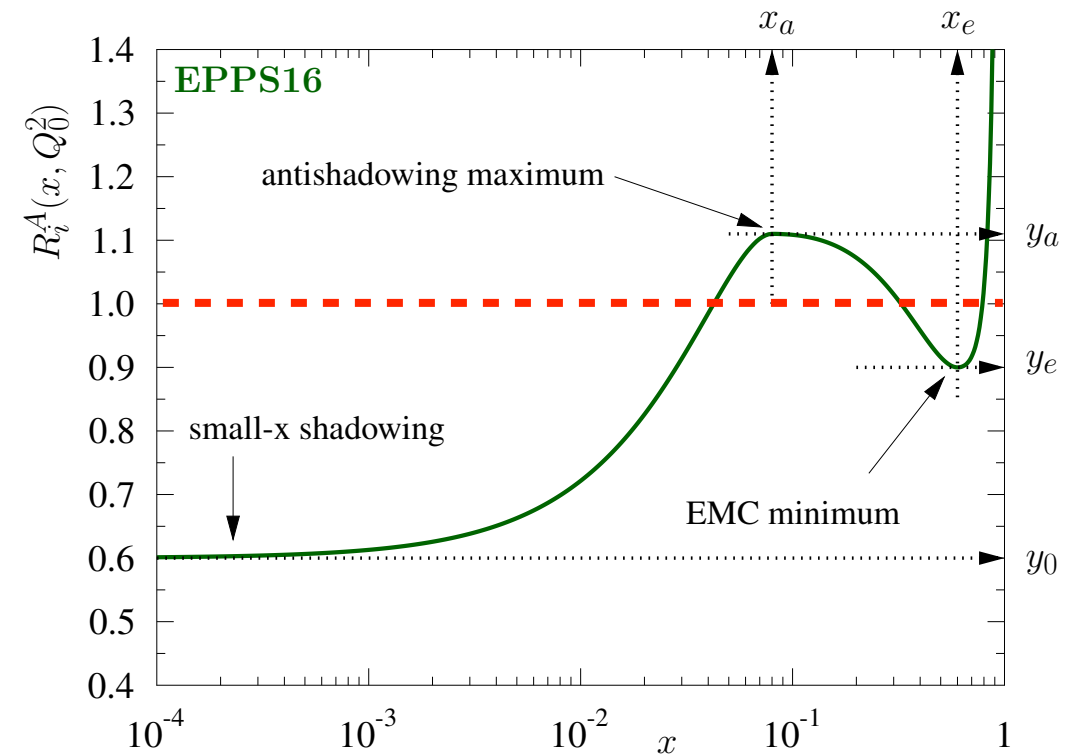
- Not being able to access the real x_1, x_2 from measuring the kinematics is not exclusive of the Drell-Yan process. It happens for all p+p collisions.

- And can have significant impact: e.g. in p+Pb collisions

$$x_1, x_2 = \sqrt{\frac{M^2}{s}} e^{\pm y}$$



data from e+A DIS



ratio of proton to nuclear PDF

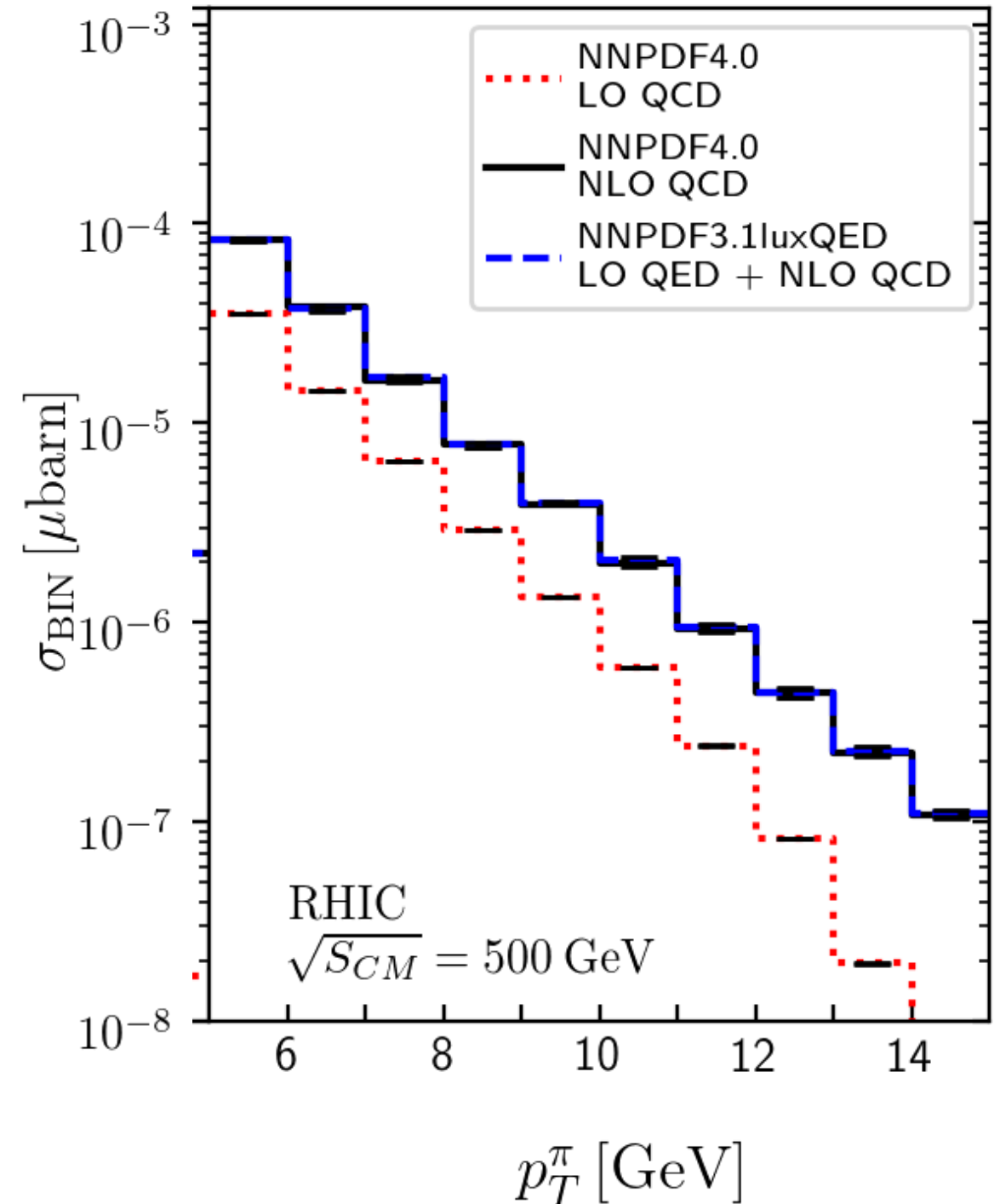
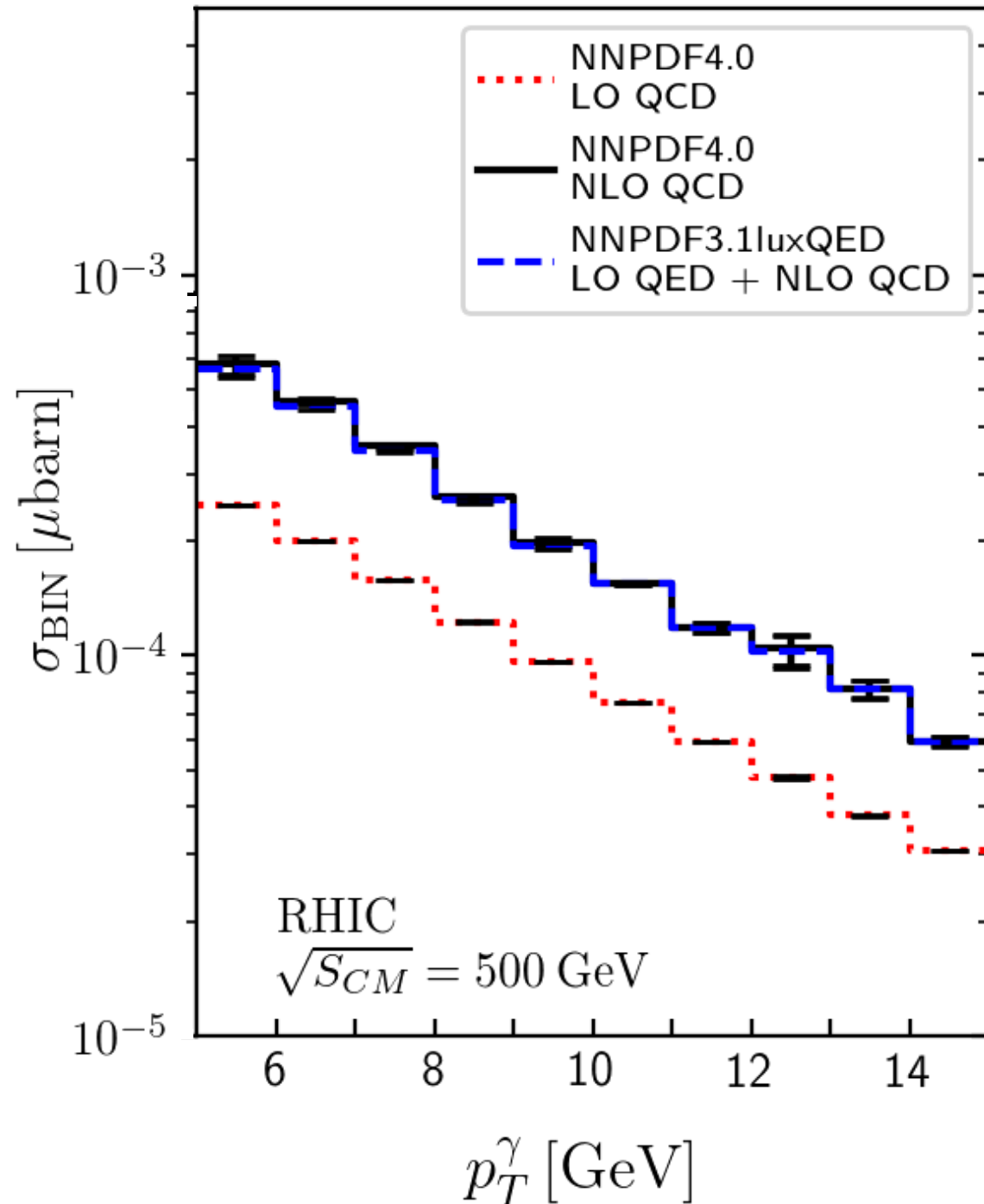
- We want to use ML to find a link between the measurable quantities and the parton momentum fractions.

- ▶ We looked at one particular process: $p + p \rightarrow \pi^+ + \gamma$
- ▶ Reconstructed x_1, x_2 and z from momenta of π^+, γ
- ▶ For RHIC kinematics, so we could compare with previous results.

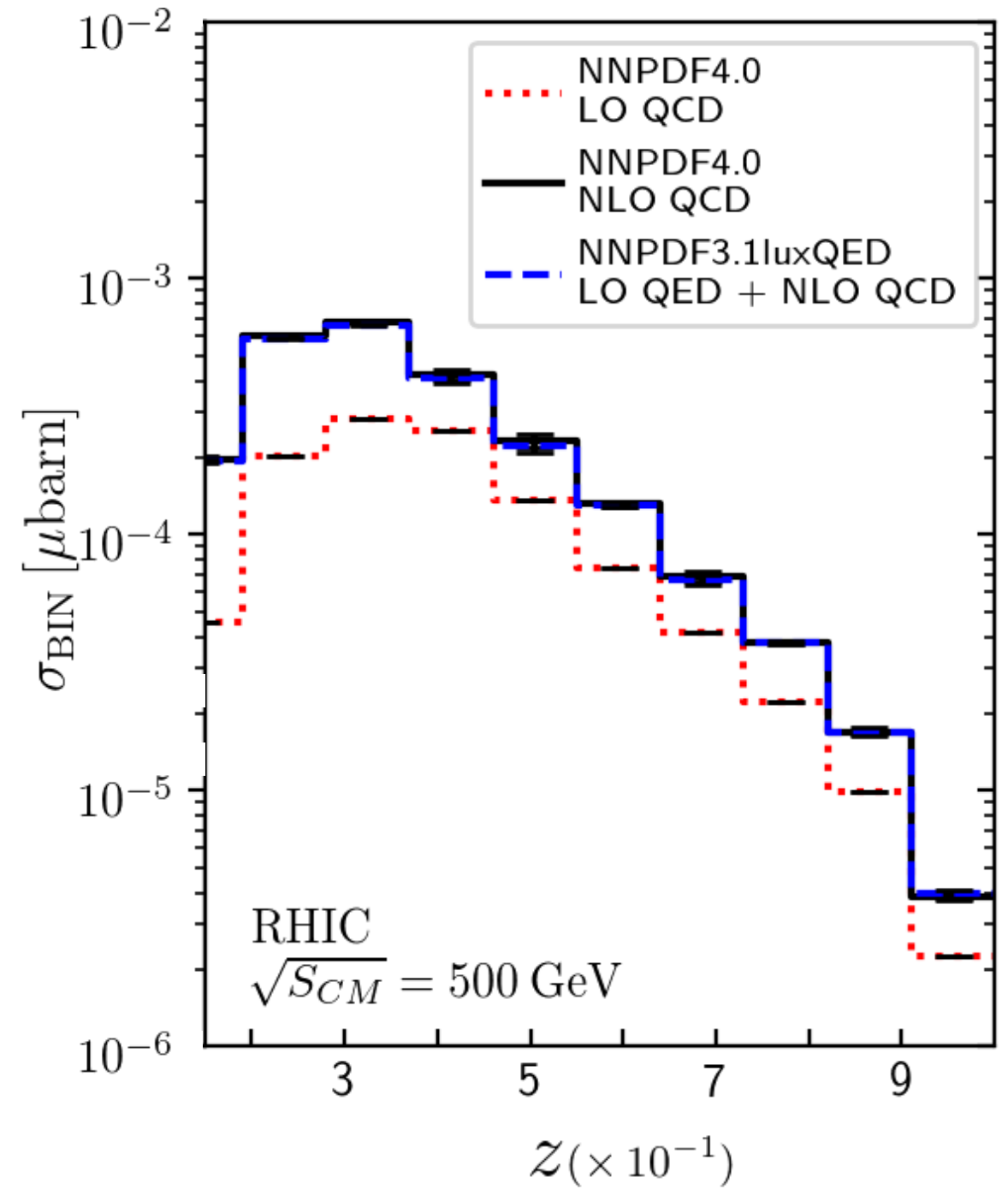
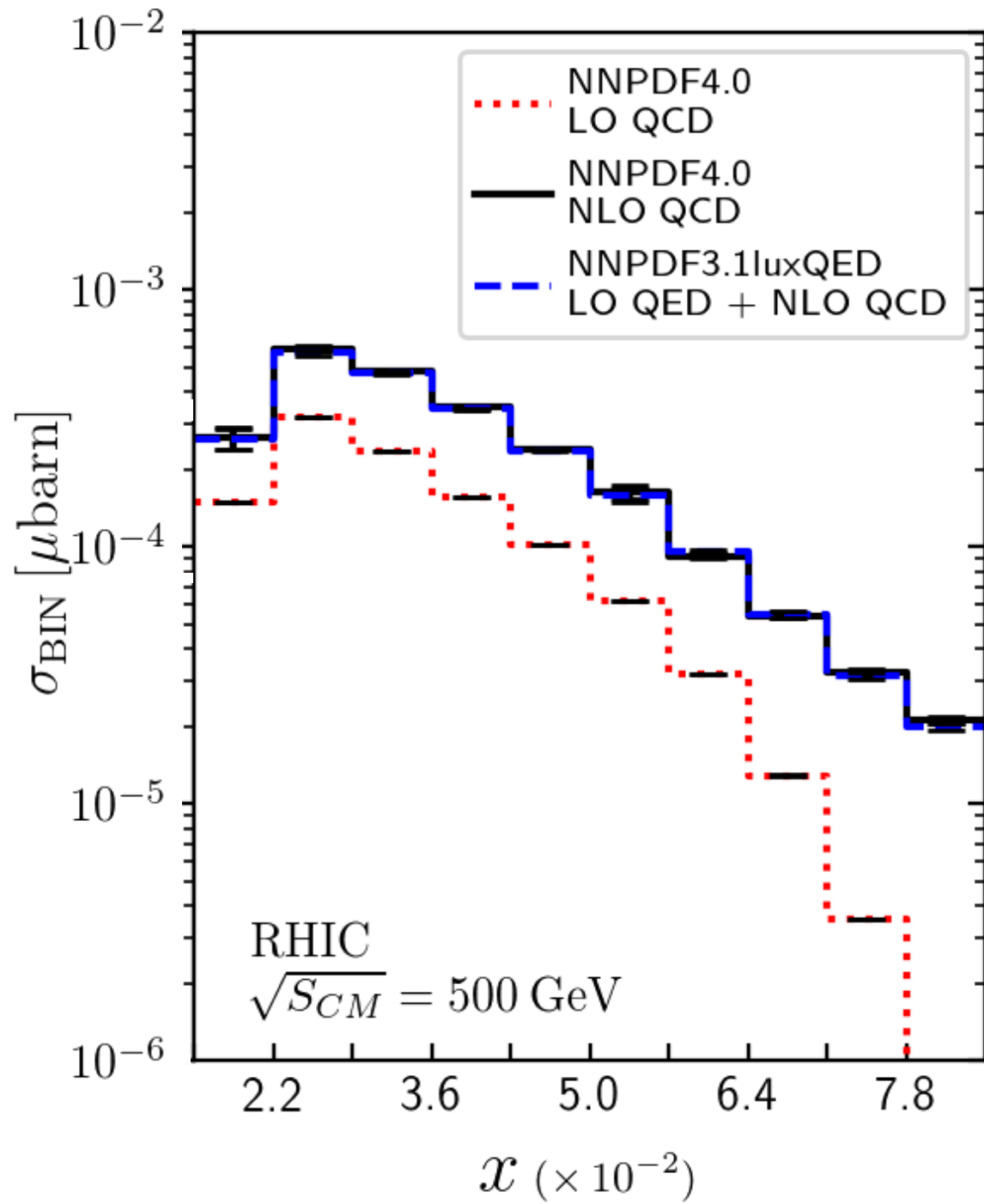
D. de Florian and G. Sborlini,
Phys.Rev.D 83 (2011) 074022

First: check the dependences on the kinematics

- ▶ Transverse momentum dependence:



- $x_1 = x, z$ dependences:



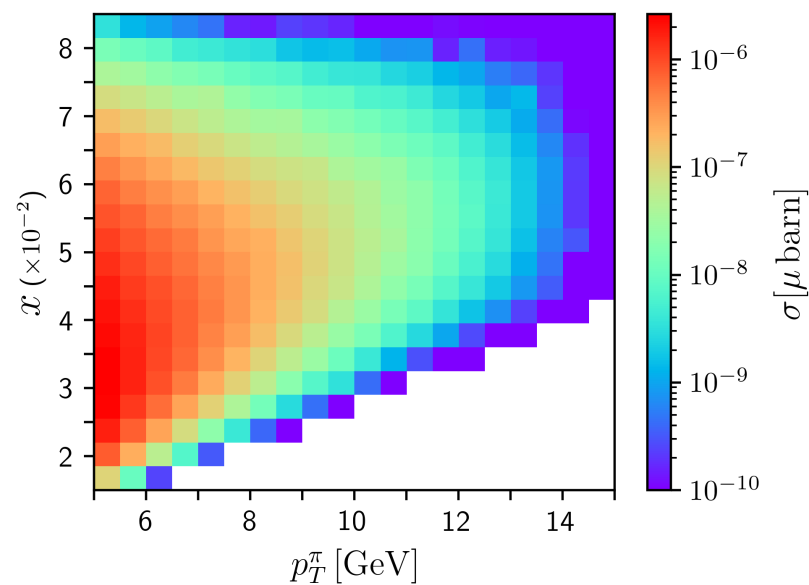
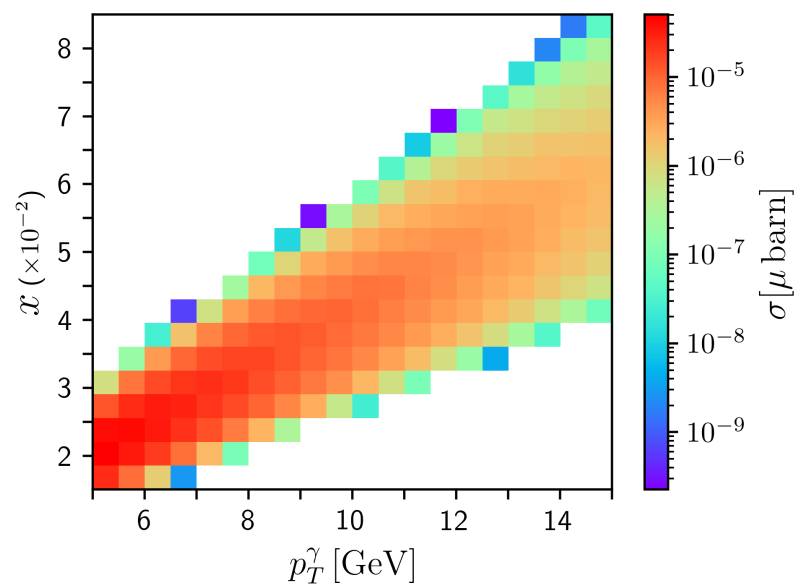
Second: check correlations

LO Kinematics

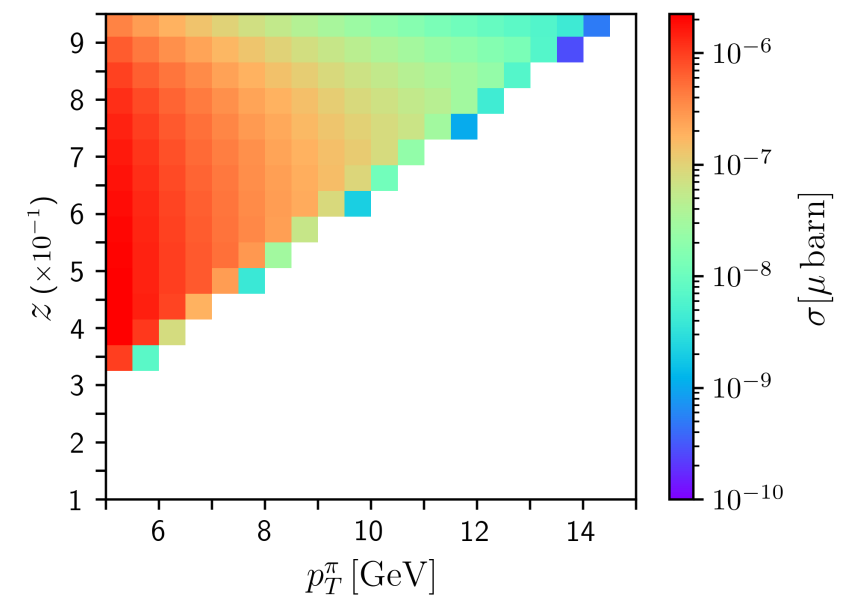
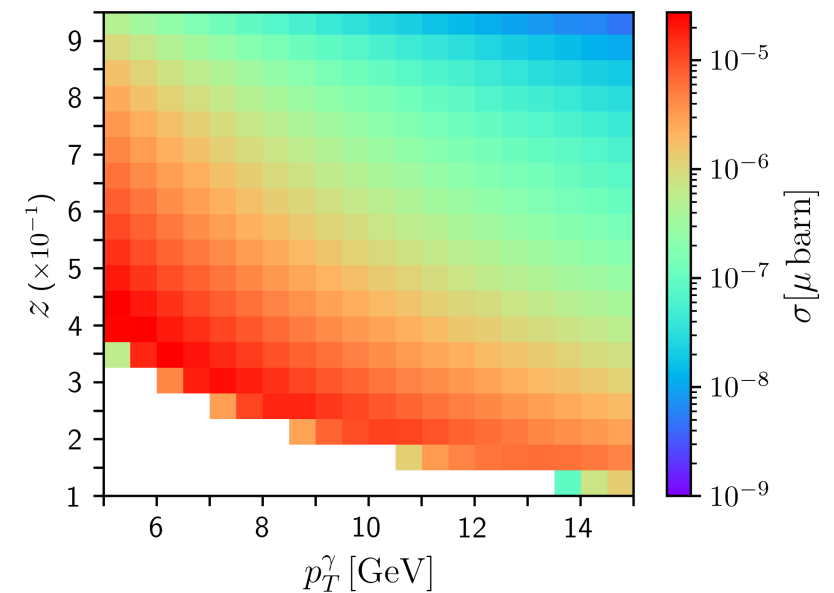
$$x_{1,2} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\eta^{\pm\pi}} + e^{\eta^{\pm\gamma}} \right)$$

$$z = \frac{p_T^\pi}{p_T^\gamma}$$

► x vs. p_T



► z vs. p_T



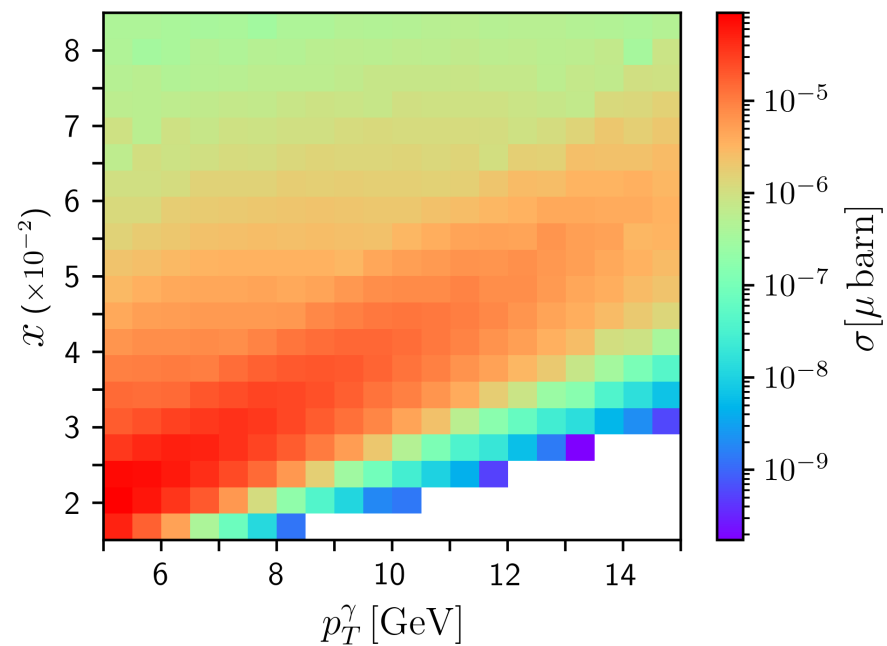
Second: check correlations

NLO Kinematics

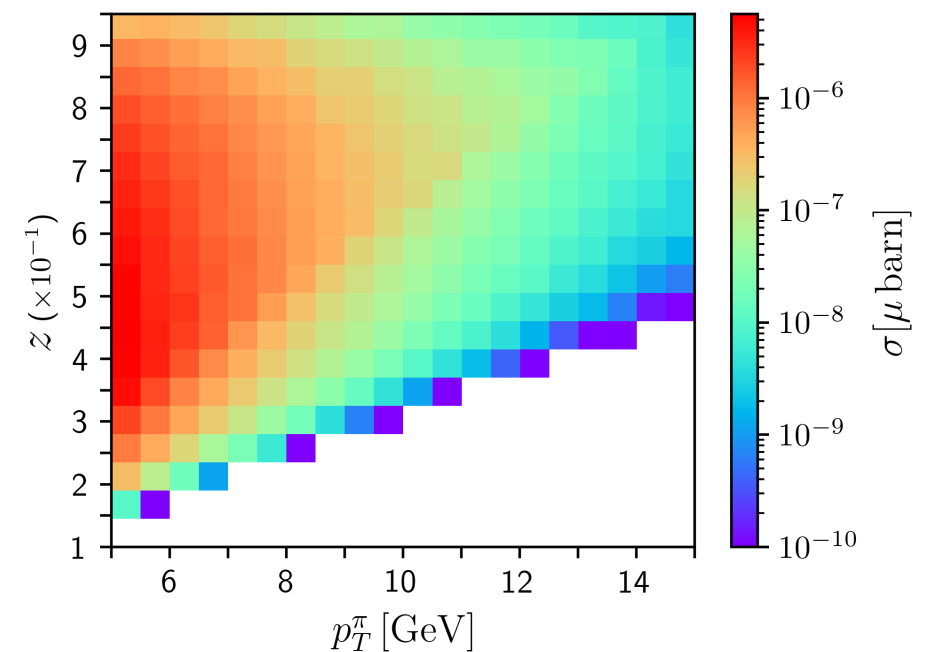
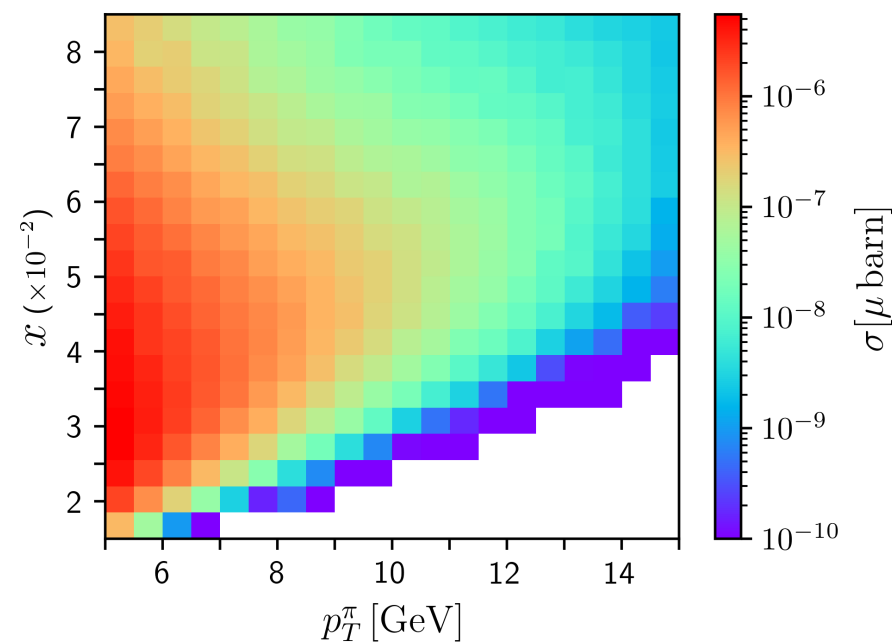
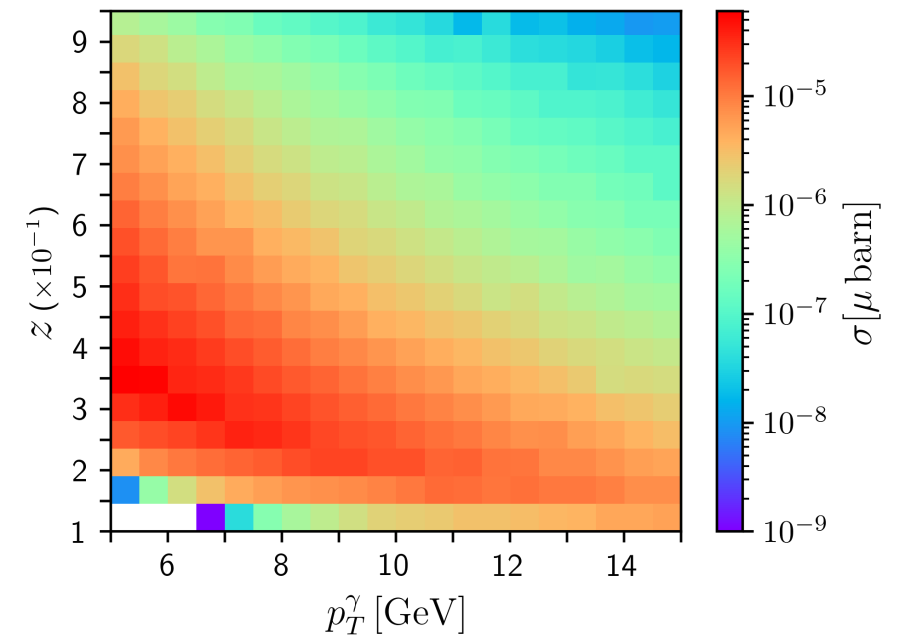
$x_{1,2} = ?$

$z = ?$

► x vs. p_T



► z vs. p_T



Kinematics: LO

$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\pm\eta^\pi} + e^{\pm\eta^\gamma} \right)$$

Kinematics: NLO

$$x_{1,2}^{rec.} = ?$$

Kinematics: LO

$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\pm\eta^\pi} + e^{\pm\eta^\gamma} \right)$$

Kinematics: NLO

$$x_{1,2}^{rec.} = ?$$

- ▶ Experimental collaborations used

$$x_{1,2}^{rec.} = \frac{p_T^\gamma e^{\pm\eta^\pi} - \cos(\phi^\pi - \phi^\gamma) p_T^\gamma e^{\pm\eta^\gamma}}{\sqrt{s}}$$

Kinematics: LO

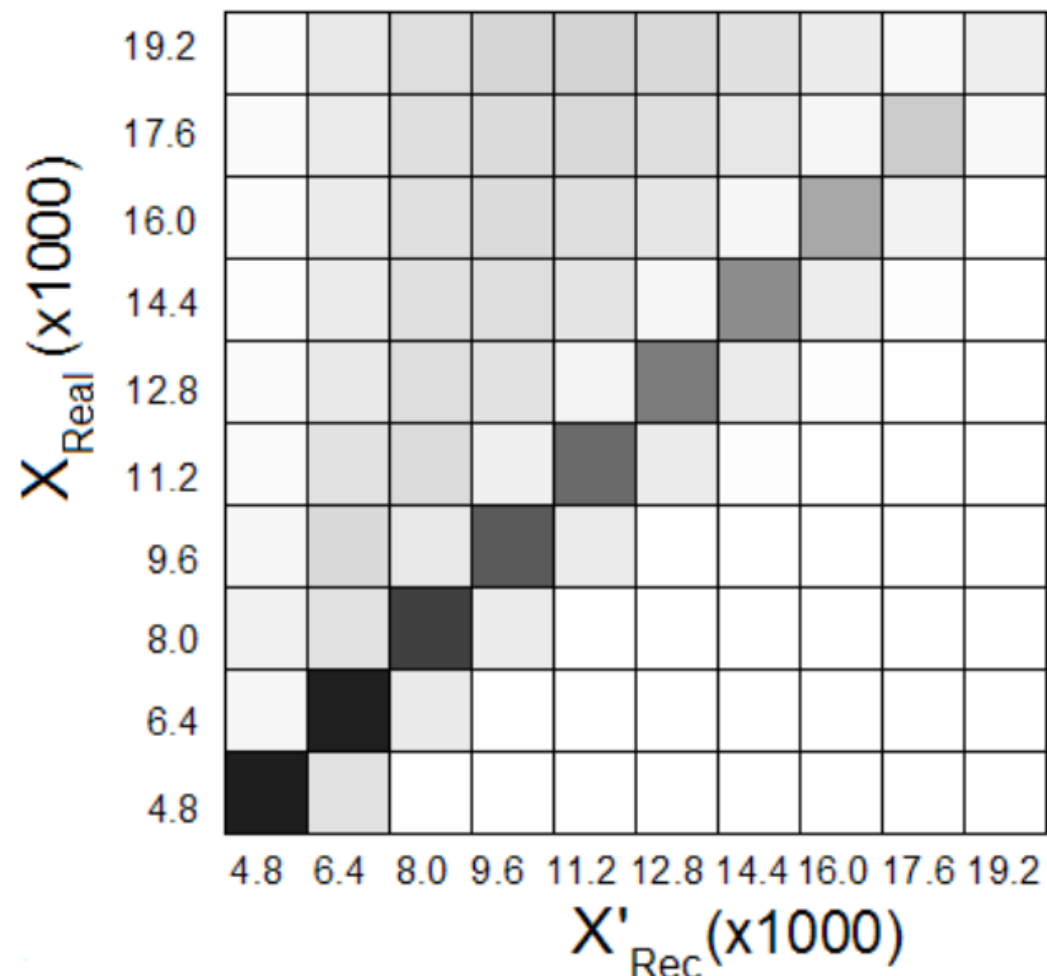
$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\pm\eta^\pi} + e^{\pm\eta^\gamma} \right)$$

Kinematics: NLO

$$x_{1,2}^{rec.} = ?$$

- Experimental collaborations used

$$x_{1,2}^{rec.} = \frac{p_T^\gamma e^{\pm\eta^\pi} - \cos(\phi^\pi - \phi^\gamma) p_T^\gamma e^{\pm\eta^\gamma}}{\sqrt{s}}$$



Kinematics: LO

$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\pm\eta^\pi} + e^{\pm\eta^\gamma} \right)$$

$$z^{rec.} = \frac{p_T^\pi}{p_T^\gamma}$$

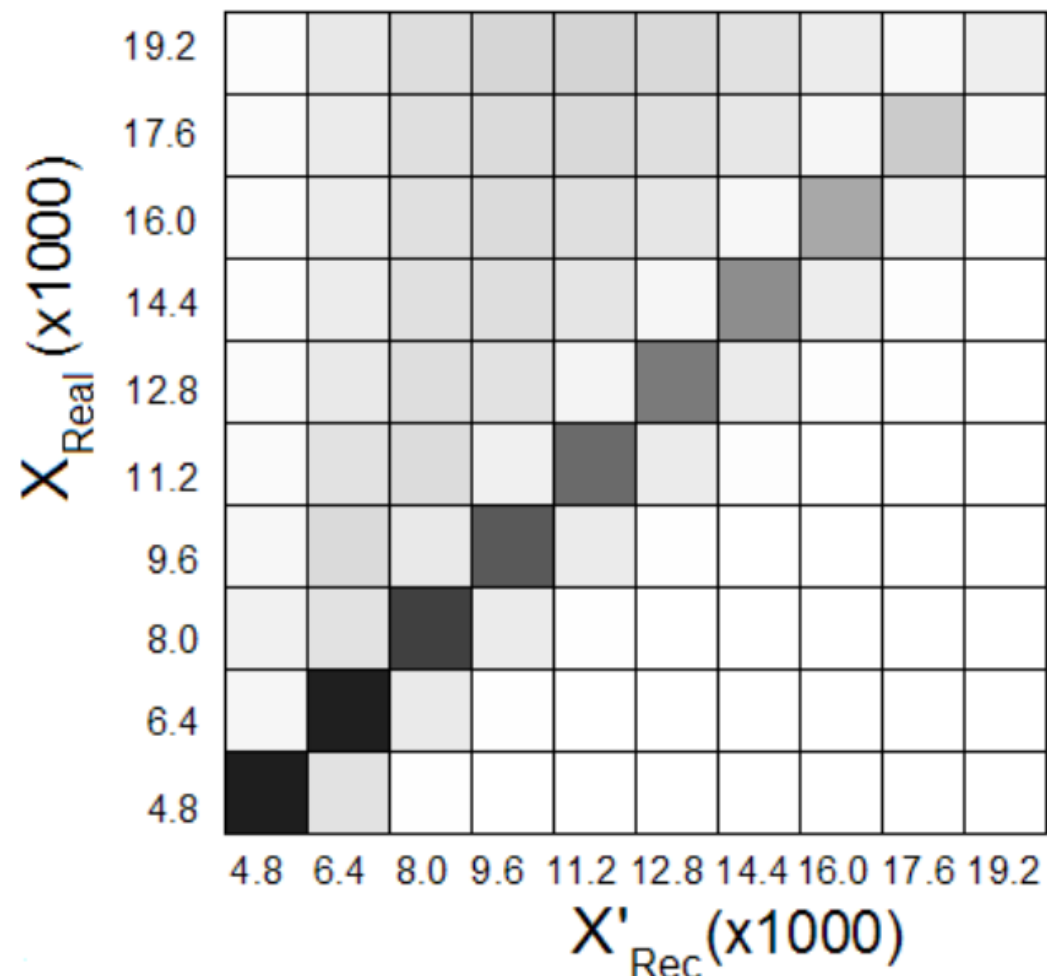
Kinematics: NLO

$$x_{1,2}^{rec.} = ?$$

$$z^{rec.} = ?$$

- Experimental collaborations used

$$z^{rec.} = -\cos(\phi^\pi - \phi^\gamma) \frac{p_T^\pi}{p_T^\gamma}$$



Kinematics: LO

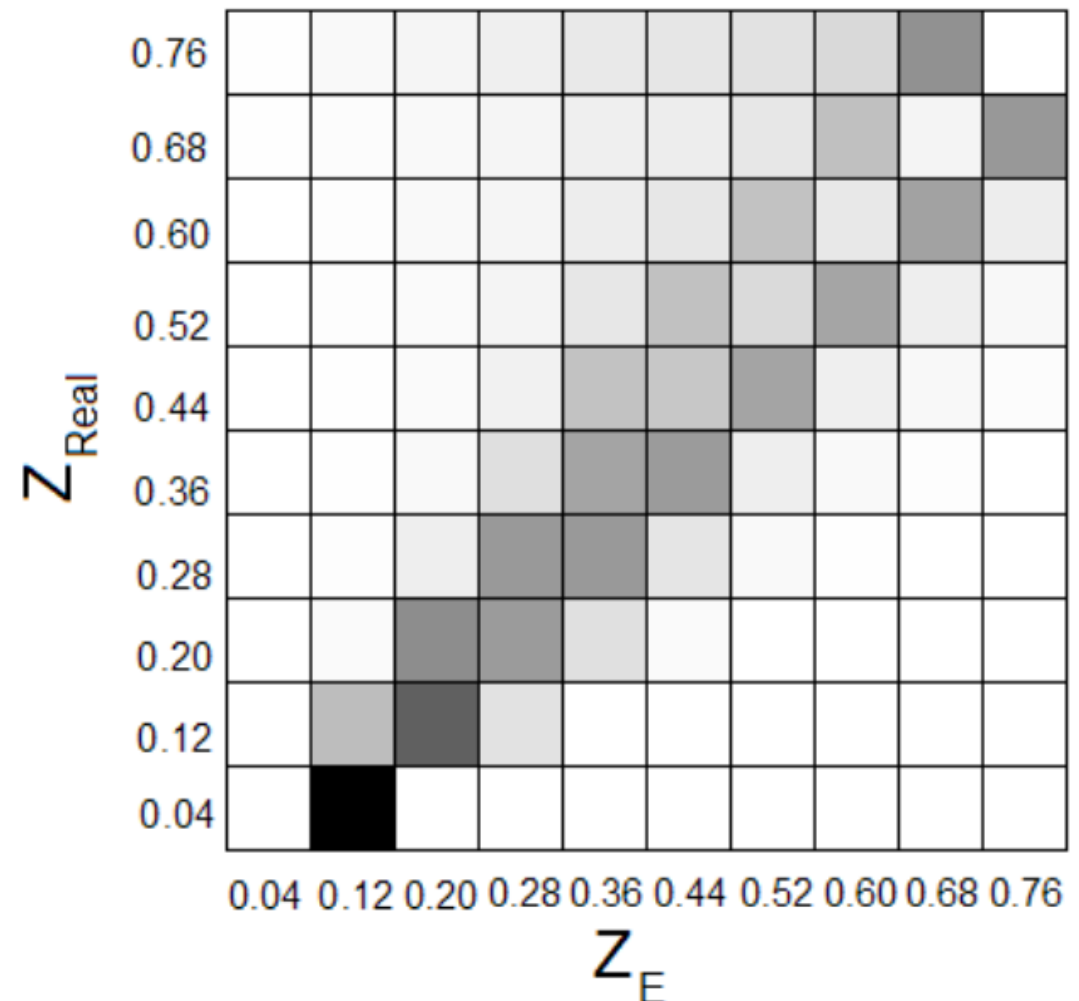
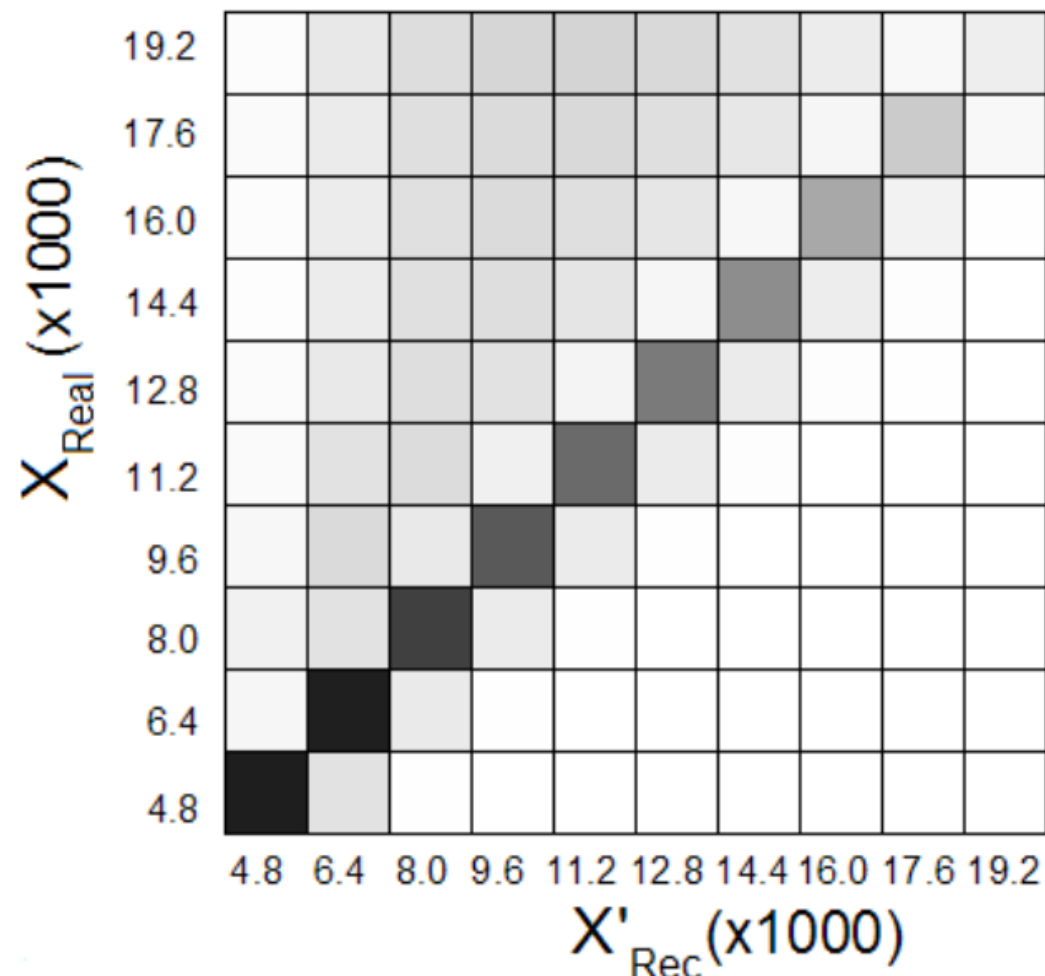
$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\pm\eta^\pi} + e^{\pm\eta^\gamma} \right) \quad z^{rec.} = \frac{p_T^\pi}{p_T^\gamma}$$

Kinematics: NLO

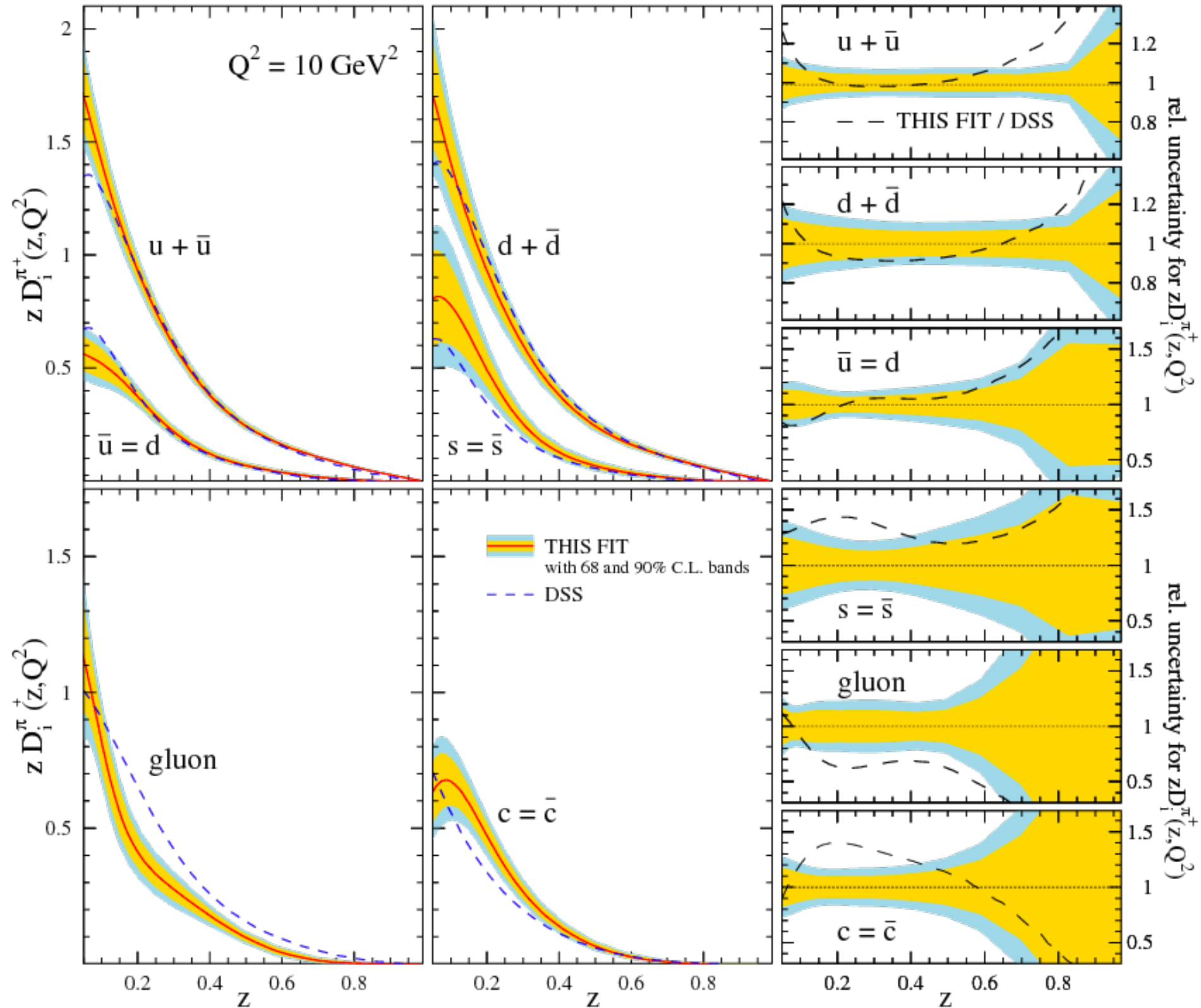
$$x_{1,2}^{rec.} = ? \quad z^{rec.} = ?$$

- Experimental collaborations used

$$z^{rec.} = -\cos(\phi^\pi - \phi^\gamma) \frac{p_T^\pi}{p_T^\gamma}$$



- In the last decade PDFs and FFs have changed significantly.



- ▶ In the last decade PDFs and FFs have changed significantly.

- ▶ The code now includes QED corrections:

D. Rentería-Estrada, R. Hernández-Pinto, G. Sborlini, *Symmetry* 13 (2021) 6, 942

- ▶ New numerical methods/tools available with tutorials “for dummies”.

- ▶ In the last decade PDFs and FFs have changed significantly.

- ▶ The code now includes QED corrections:

D. Rentería-Estrada, R. Hernández-Pinto, G. Sborlini, *Symmetry* 13 (2021) 6, 942

- ▶ New numerical methods/tools available with tutorials “for dummies”.
- ▶ We’re the dummies: we want to apply machine learning techniques to access the real momentum fractions and lower the assumptions.

- ▶ At NLO we have real ($2 \rightarrow 3$) and virtual ($2 \rightarrow 2$) contributions and counterterms ($2 \rightarrow 2$).
- ▶ Cancellations can only happen in the MC integration when histogramming.

$$\{\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)\} \in \bar{\mathcal{V}}_{EXP}$$

$$\sigma_j(\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)) = \int_{(p_T^\gamma)_{j,MIN}}^{(p_T^\gamma)_{j,MAX}} dp_T^\gamma \int_{(p_T^\pi)_{j,MIN}}^{(p_T^\pi)_{j,MAX}} dp_T^\pi \int dx_1 dx_2 dz d\bar{\sigma}$$

- ▶ At NLO we have real ($2 \rightarrow 3$) and virtual ($2 \rightarrow 2$) contributions and counterterms ($2 \rightarrow 2$).
- ▶ Cancellations can only happen in the MC integration when histogramming.

$$\{\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)\} \in \bar{\mathcal{V}}_{EXP}$$

$$\sigma_j(\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)) = \int_{(p_T^\gamma)_{j,MIN}}^{(p_T^\gamma)_{j,MAX}} dp_T^\gamma \int_{(p_T^\pi)_{j,MIN}}^{(p_T^\pi)_{j,MAX}} dp_T^\pi \int dx_1 dx_2 dz d\bar{\sigma}$$

- ▶ We weight the momentum fractions from the MC with the per-bin cross-section

$$(x_1)_j = \sum_i (x_1)_i \frac{d\sigma_j}{dx_1}(p_j; (x_1)_i)$$

$$(z)_j = \sum_i (z)_i \frac{d\sigma_j}{dz}(p_j; (z)_i)$$

- ▶ At NLO we have real ($2 \rightarrow 3$) and virtual ($2 \rightarrow 2$) contributions and counterterms ($2 \rightarrow 2$).
- ▶ Cancellations can only happen in the MC integration when histogramming.

$$\{\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)\} \in \bar{\mathcal{V}}_{EXP}$$

$$\sigma_j(\bar{p}_T^\gamma, \bar{p}_T^\pi, \bar{\eta}^\gamma, \bar{\eta}^\pi, \overline{\cos}(\phi^\pi - \phi^\gamma)) = \int_{(p_T^\gamma)_{j,MIN}}^{(p_T^\gamma)_{j,MAX}} dp_T^\gamma \int_{(p_T^\pi)_{j,MIN}}^{(p_T^\pi)_{j,MAX}} dp_T^\pi \int dx_1 dx_2 dz d\bar{\sigma}$$

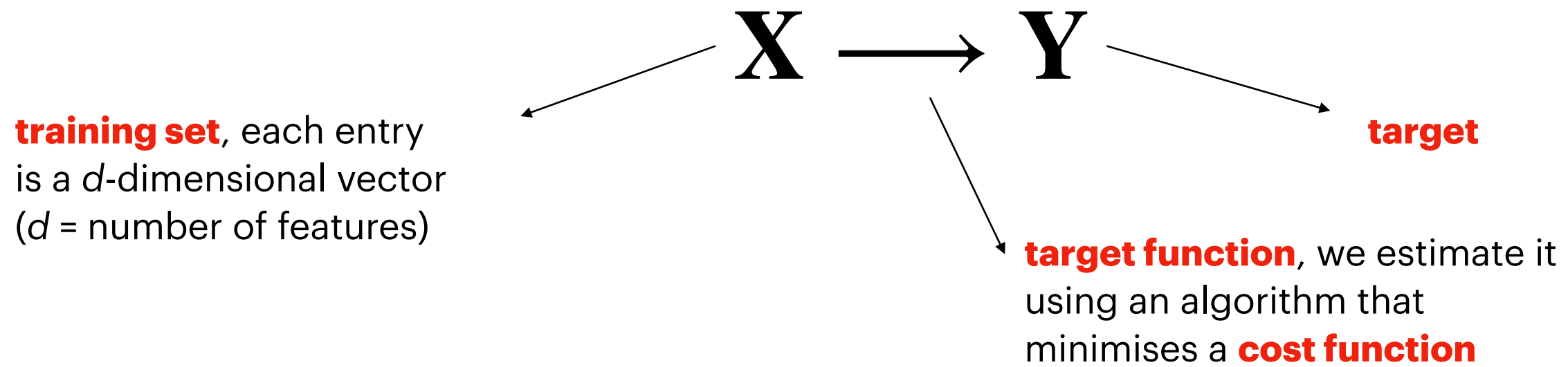
- ▶ We weight the momentum fractions from the MC with the per-bin cross-section

$$(x_1)_j = \sum_i (x_1)_i \frac{d\sigma_j}{dx_1}(p_j; (x_1)_i) \qquad (z)_j = \sum_i (z)_i \frac{d\sigma_j}{dz}(p_j; (z)_i)$$

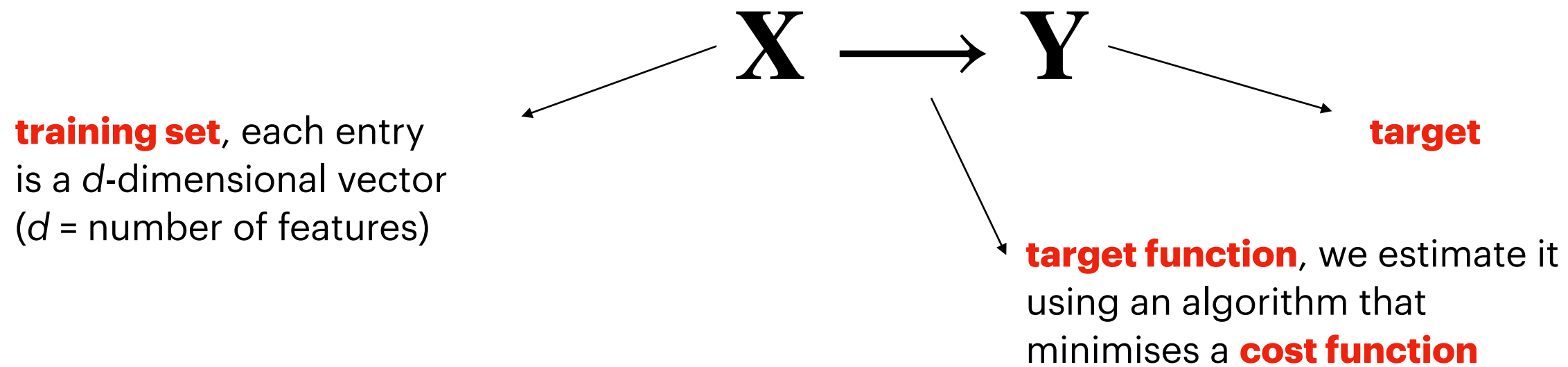
- ▶ With this we search for the mapping

$$X_{1,REC} : \bar{\mathcal{V}}_{EXP} \rightarrow \bar{X}_{1,REAL} = \{(x_1)_j\}$$

- ▶ In general, in ML



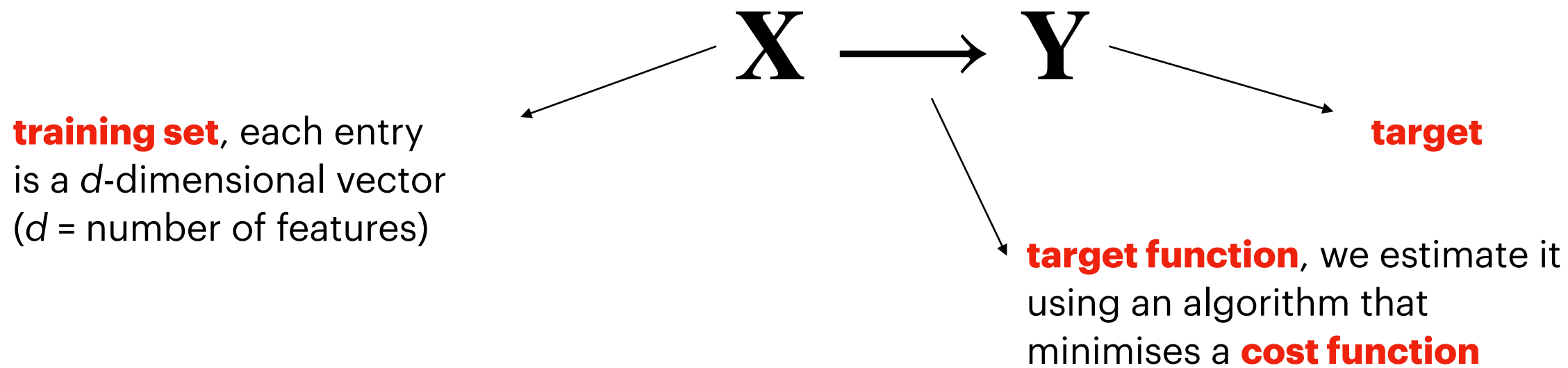
- ▶ In general, in ML



Linear regression

$$\hat{y} = \theta_0 + \theta_1 x^{(1)}$$

- ▶ In general, in ML



Linear regression

$$\hat{y} = \theta_0 + \theta_1 x^{(1)}$$

- ▶ “cheat”: **linear** means linear *in the parameters*.

$$[x^{(1)}]^2 = x^{(2)} \longrightarrow \hat{y} = \theta_0 + \theta_1 x^{(1)} + \theta_2 x^{(2)}$$

- ▶ the parameters minimise

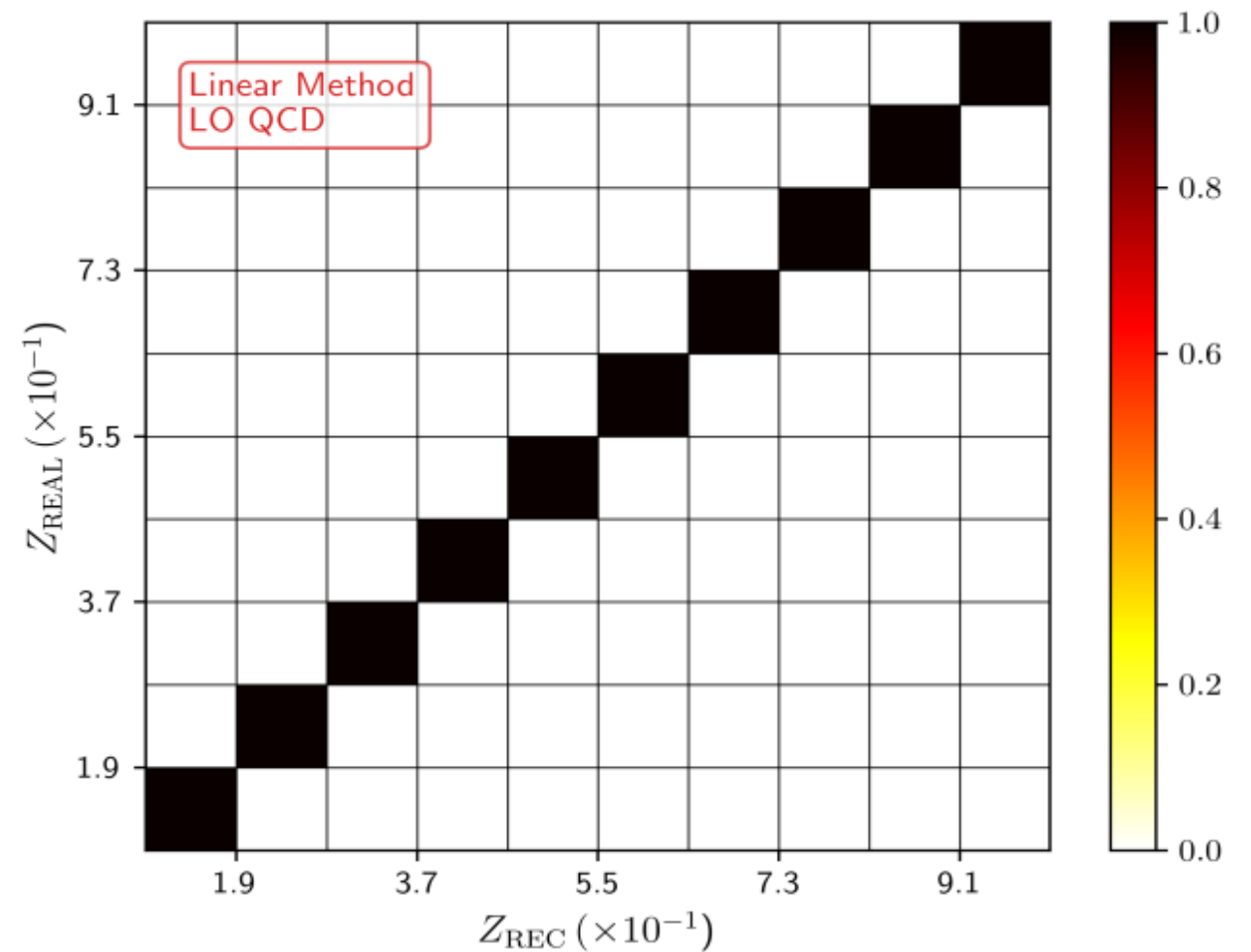
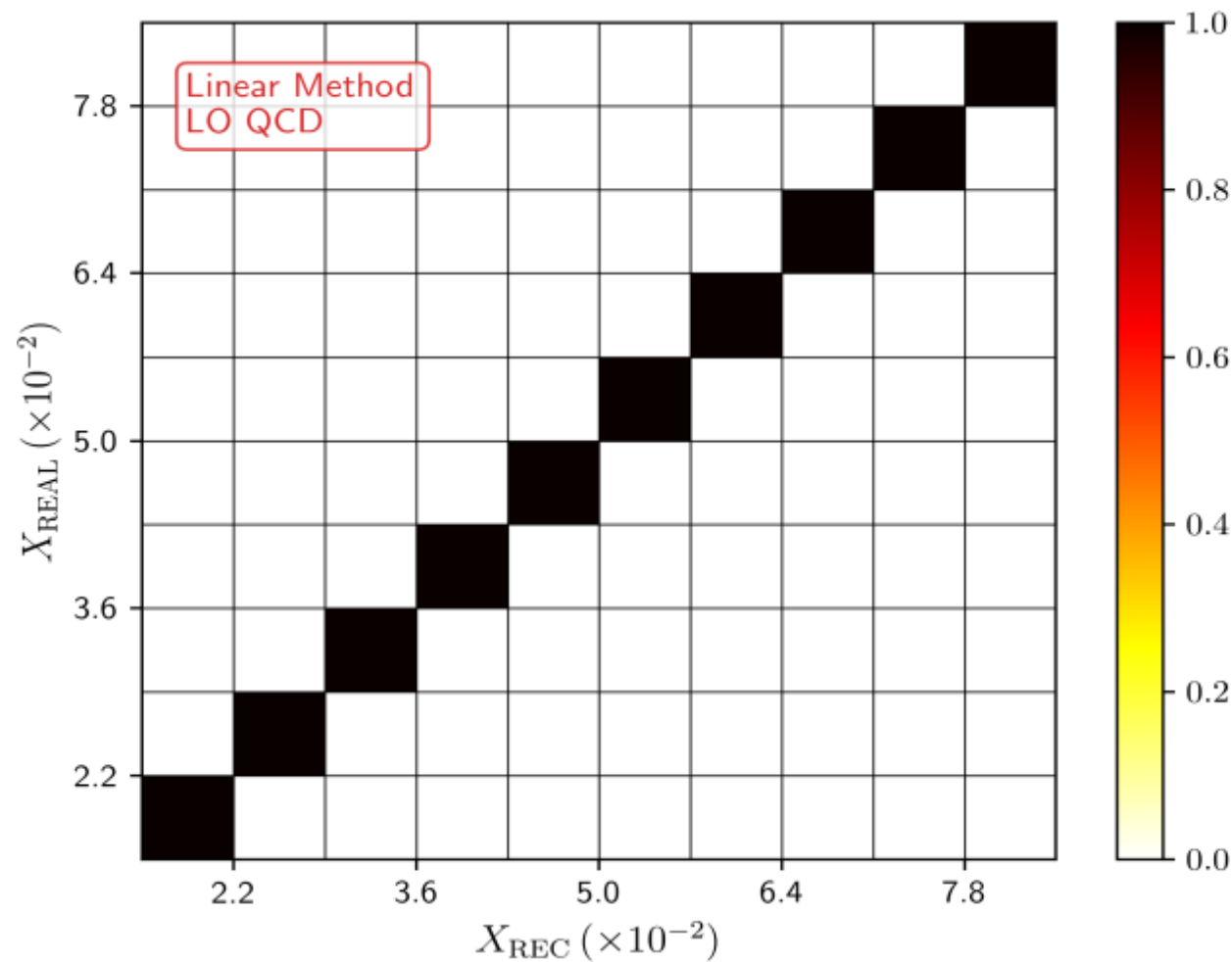
$$\min_{\theta} \|\hat{y} - y\|_2^2$$

- ▶ Let us start with LO and use linear regression:

▶ Basis: $\mathcal{B}_{LO} = \left\{ \frac{p_T^\gamma e^{\eta^\pi}}{\sqrt{s}}, \frac{p_T^\gamma e^{\eta^\gamma}}{\sqrt{s}}, \frac{p_T^\gamma e^{-\eta^\pi}}{\sqrt{s}}, \frac{p_T^\gamma e^{-\eta^\gamma}}{\sqrt{s}}, \frac{p_T^\pi}{p_T^\gamma} \right\}$

$$x_{1,2}^{rec.} = \frac{p_T^\gamma}{\sqrt{s}} \left(e^{\eta^{\pm\pi}} + e^{\eta^{\pm\gamma}} \right)$$

$$z^{rec.} = \frac{p_T^\pi}{p_T^\gamma}$$



Linear regression

- ▶ We used three bases: “LO inspired”, “general”, “physically motivated”

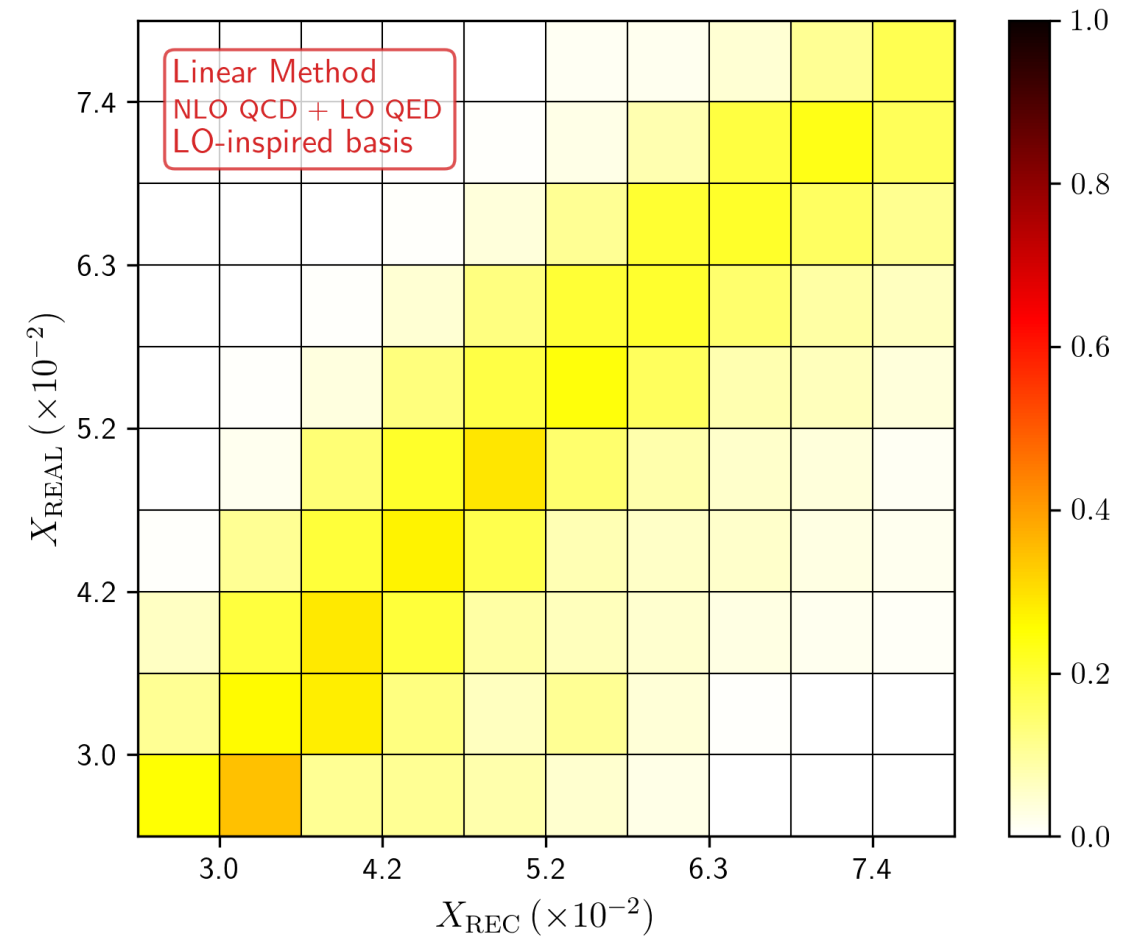
$$x_1^{LO} = \frac{p_T^\gamma e^{\eta^\pi} + p_T^\gamma e^{\eta^\gamma}}{\sqrt{s}}$$

$$x_1^{prop.} = \frac{p_T^\gamma e^{\eta^\pi} - \cos(\phi^\pi - \phi^\gamma) p_T^\gamma e^{\eta^\gamma}}{\sqrt{s}}$$

$$\mathcal{B}_{LO} = \left\{ \frac{p_T^\gamma e^{\eta^\pi}}{\sqrt{s}}, \frac{p_T^\gamma e^{\eta^\gamma}}{\sqrt{s}}, \frac{p_T^\gamma e^{-\eta^\pi}}{\sqrt{s}}, \frac{p_T^\gamma e^{-\eta^\gamma}}{\sqrt{s}}, \frac{p_T^\pi}{p_T^\gamma} \right\}$$

$$\mathcal{B}_{NLO}^{X_1} = \left\{ \frac{p_T^\gamma}{\sqrt{S_{CM}}} \exp(\eta^\gamma), \frac{p_T^\gamma}{\sqrt{S_{CM}}} \exp(\eta^\pi), \frac{p_T^\pi}{\sqrt{S_{CM}}} \exp(\eta^\gamma), \frac{p_T^\pi}{\sqrt{S_{CM}}} \exp(\eta^\pi), \frac{p_T^\gamma \mathcal{K}_5}{\sqrt{S_{CM}}} \exp(\eta^\gamma), \frac{p_T^\gamma \mathcal{K}_5}{\sqrt{S_{CM}}} \exp(\eta^\pi), \frac{p_T^\pi \mathcal{K}_5}{\sqrt{S_{CM}}} \exp(\eta^\gamma), \frac{p_T^\pi \mathcal{K}_5}{\sqrt{S_{CM}}} \exp(\eta^\pi) \right\}$$

$$\mathcal{K}_5 = \cos(\phi^\pi - \phi^\gamma)$$



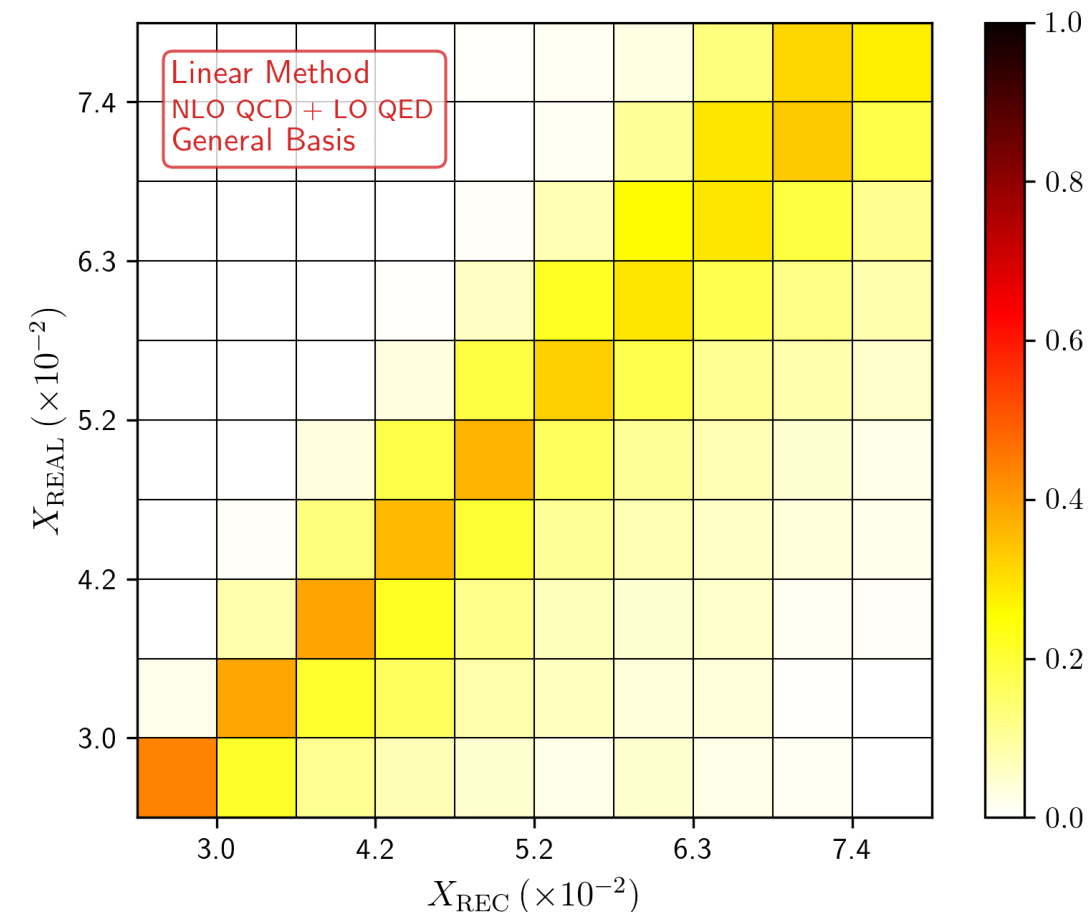
Linear regression

- We used three bases: “LO inspired”, “**general**”, “physically motivated”

$$\mathcal{K} = \left\{ \frac{p_T^\gamma}{\sqrt{s}}, \frac{p_T^\pi}{\sqrt{s}}, e^{\eta^\gamma}, e^{\eta^\pi}, \cos(\phi^\pi - \phi^\gamma) \right\}$$

$$X_{\text{REC}} = \sum_{i=1, i \neq 5}^9 (a_i + b_i \mathcal{K}_5) \mathcal{K}_i + \sum_{i \leq j, \{i,j\} \neq 5, j-i \neq 5} (c_{ij} + d_{ij} \mathcal{K}_5) \mathcal{K}_i \mathcal{K}_j$$

- 81 parameters in total.



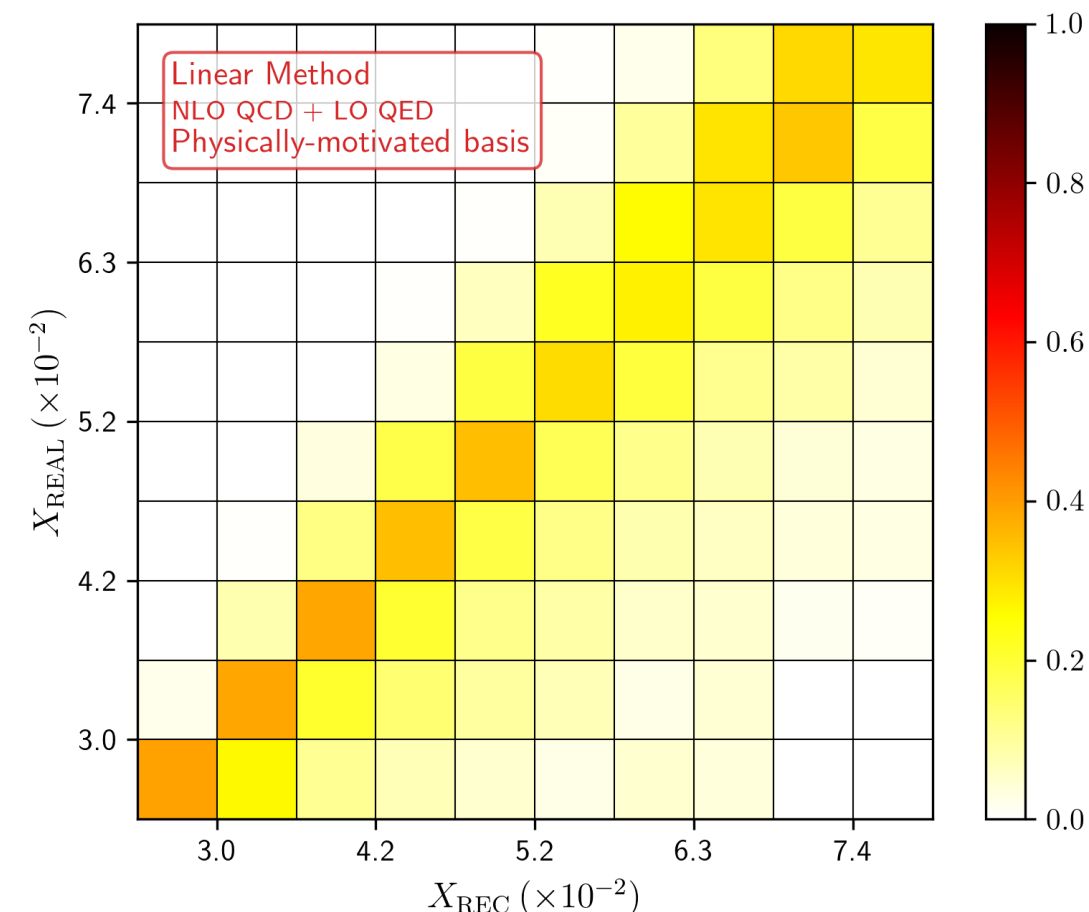
Linear regression

- We used three bases: “LO inspired”, “general”, “**physically motivated**”

$$\mathcal{K} = \left\{ \frac{p_T^\gamma}{\sqrt{s}}, \frac{p_T^\pi}{\sqrt{s}}, e^{\eta^\gamma}, e^{\eta^\pi}, \cos(\phi^\pi - \phi^\gamma) \right\}$$

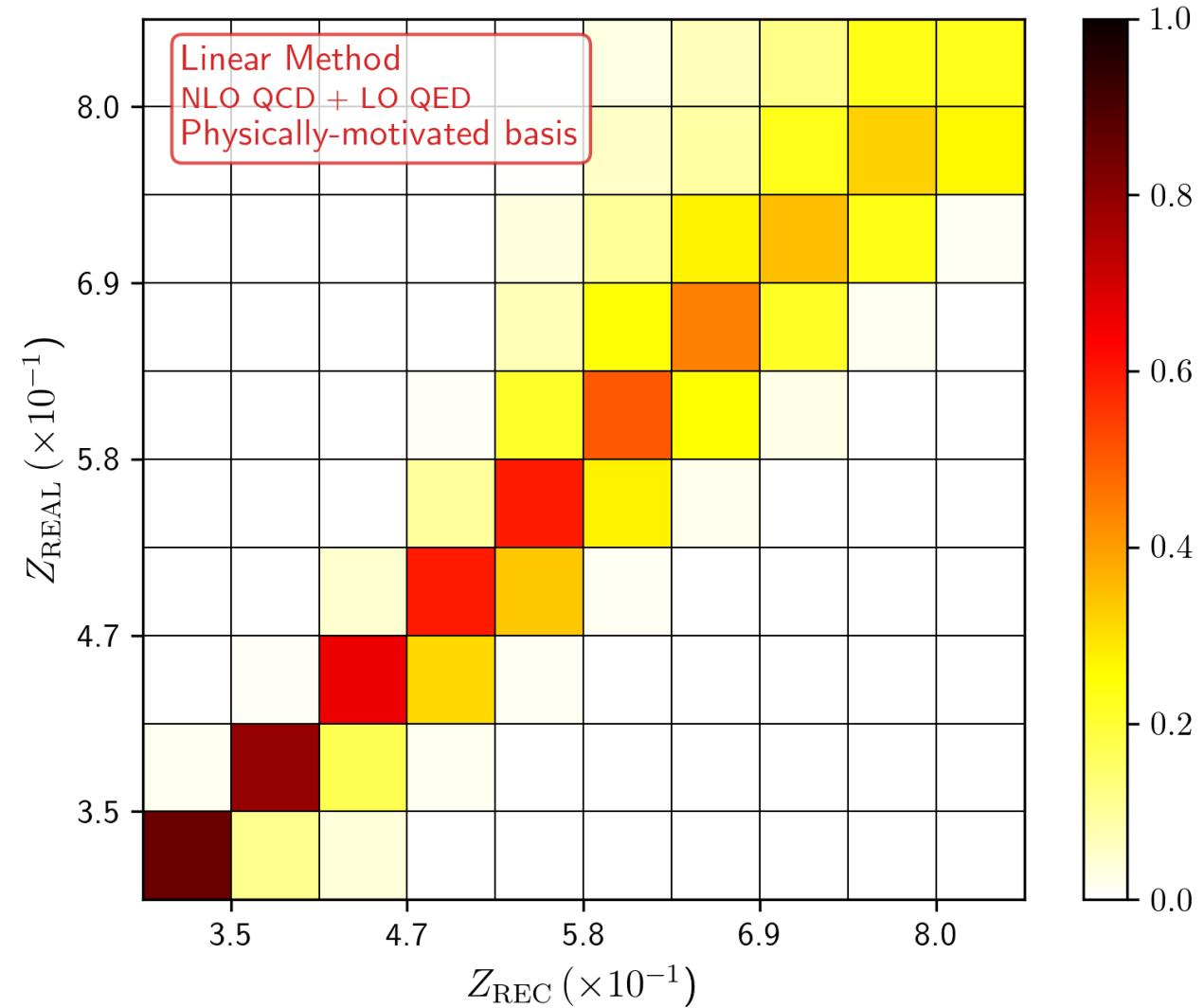
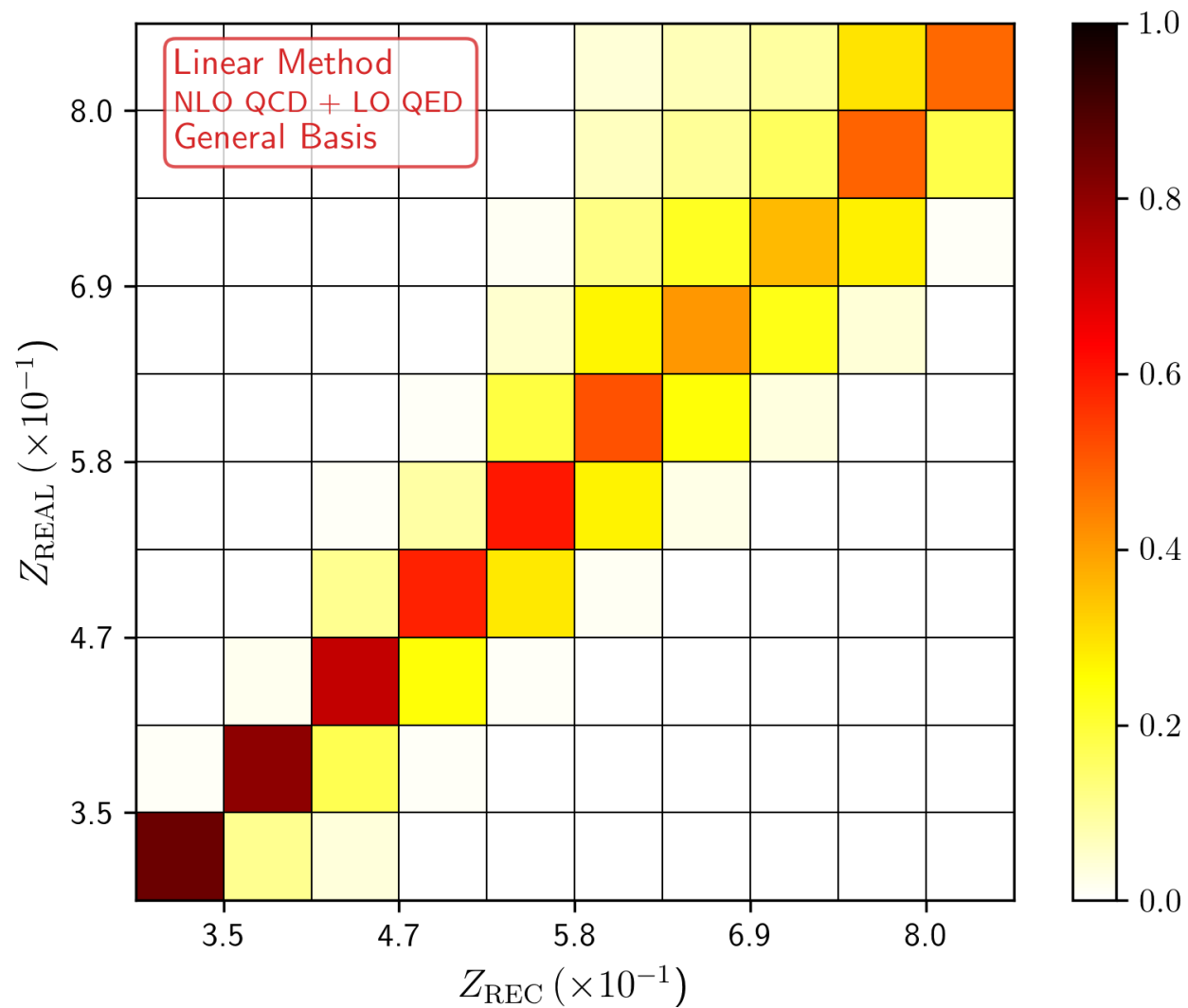
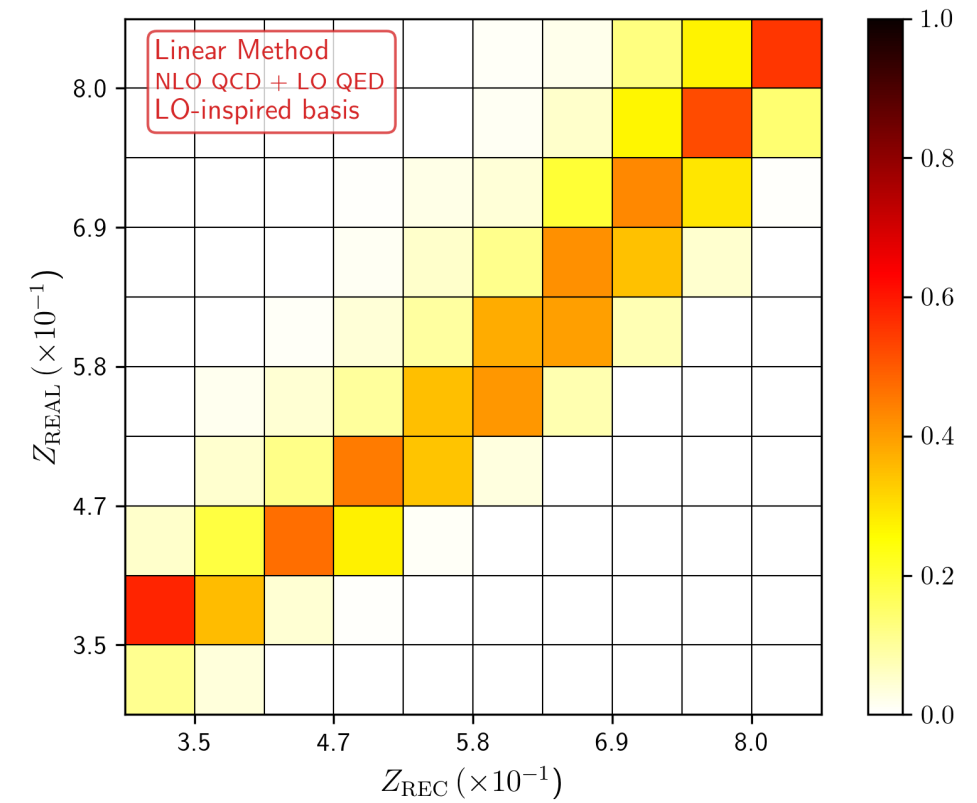
$$X_{\text{REC}} = \sum_{i=1, i \neq 5}^9 (a_i + b_i \mathcal{K}_5) \mathcal{K}_i + \sum_{i \leq j, \{i,j\} \neq 5, j-i \neq 5} (c_{ij} + d_{ij} \mathcal{K}_5) \mathcal{K}_i \mathcal{K}_j$$

- Remove contributions w.r.t. which we see no dependence (~40 parameters).



Linear regression

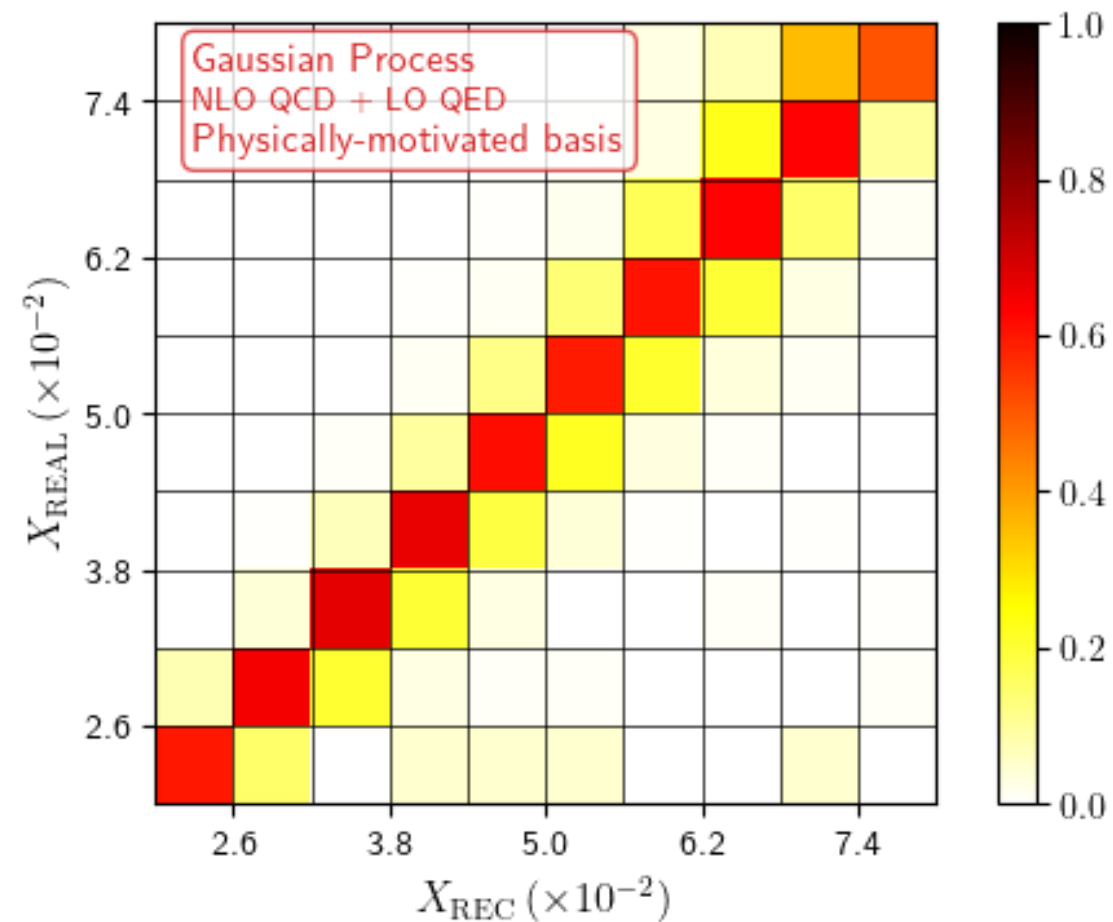
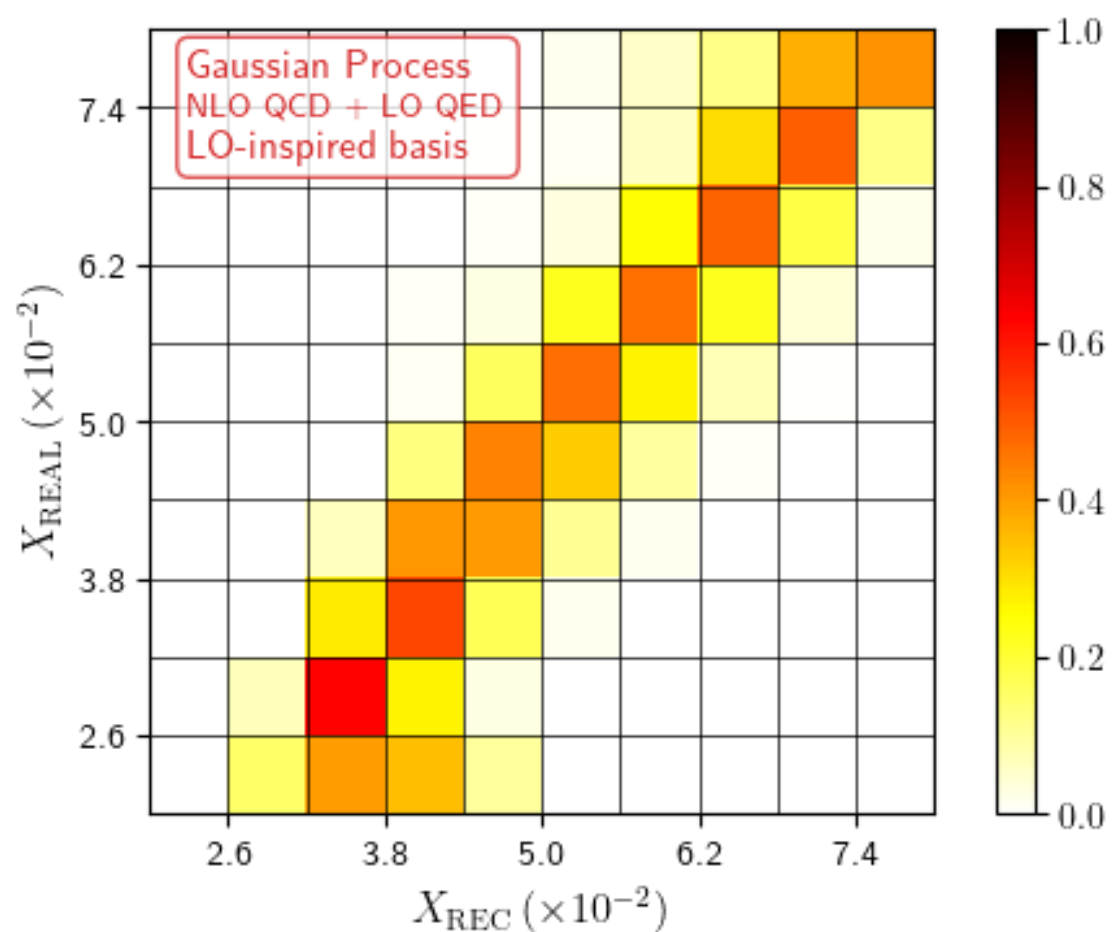
► For z:



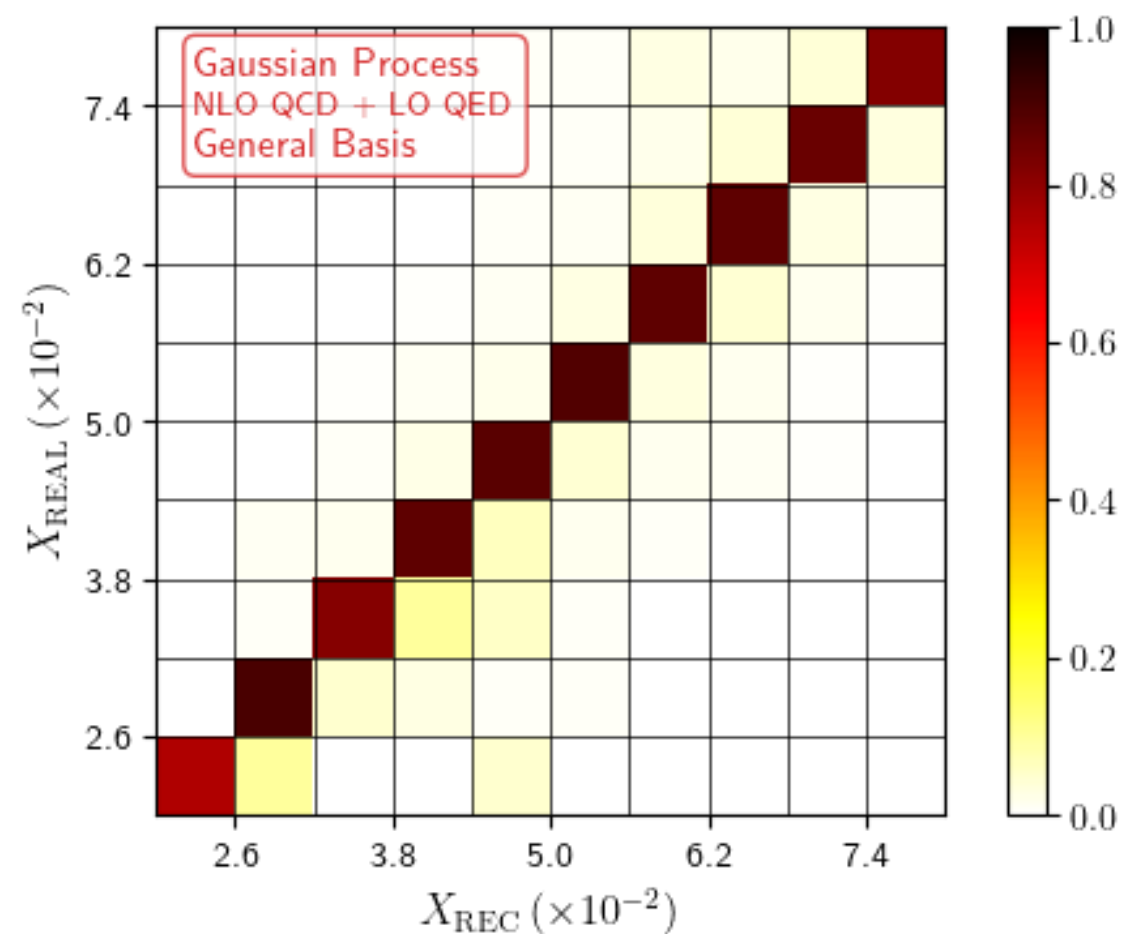
The Radial Basis Function

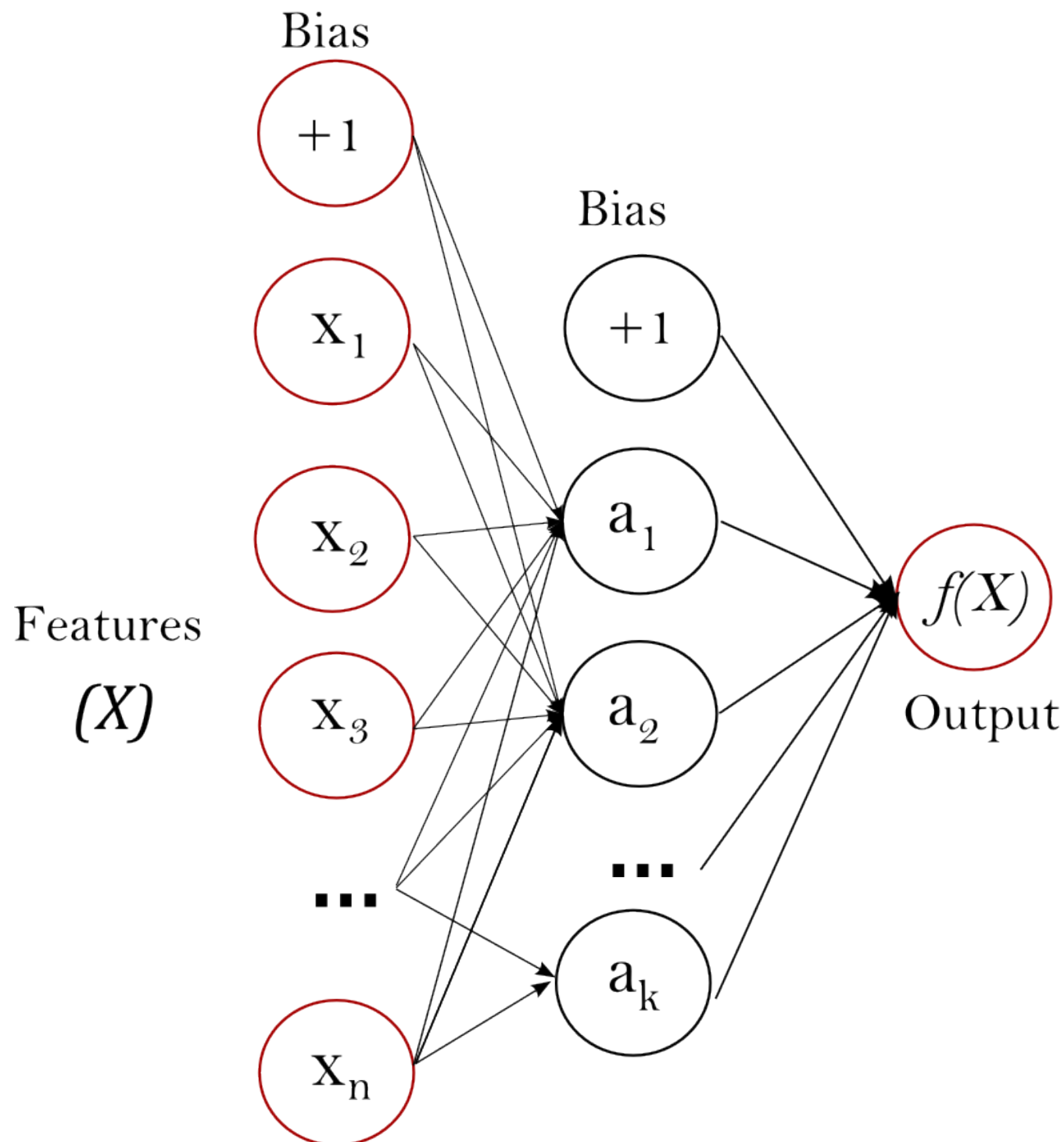
$$\mathbf{X} = (1, x_1, x_2)^T \longrightarrow \tilde{\mathbf{X}} = (1, x_1, x_2, x_1^2, x_1x_2, x_2^2)^T$$

- ▶ How to pick $\tilde{\mathbf{X}}$? Before we did it by intuition.
- ▶ We replace $x^{(i)}$ by $f_i(x) = e^{-\frac{\|x - x^{(i)}\|_2^2}{2l^2}}$
- ▶ This Radial Basis Function
 - ▶ requires less elements (one per feature in the basis).
 - ▶ effectively considers infinite terms.
 - ▶ is a popular form of the kernel method.



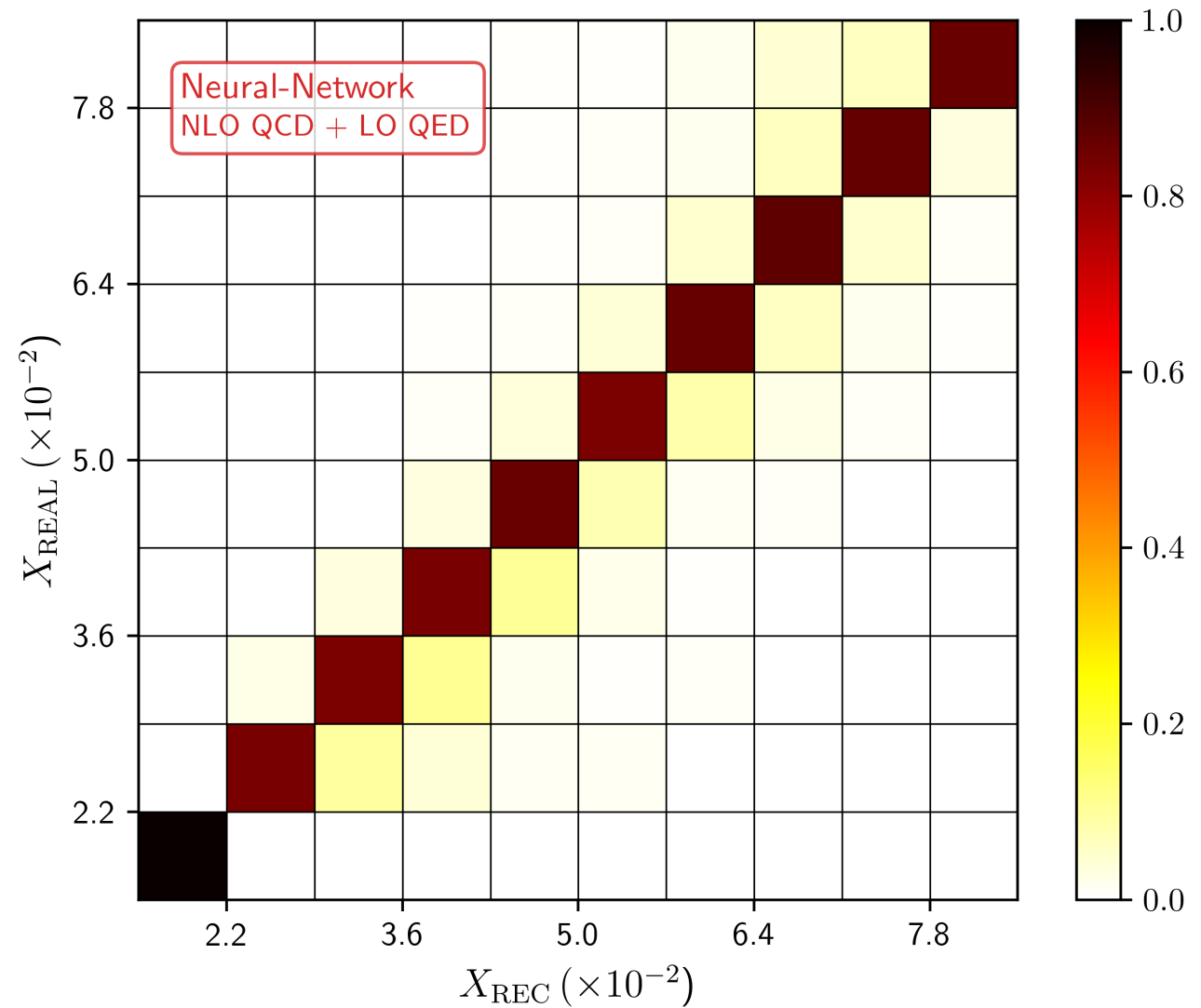
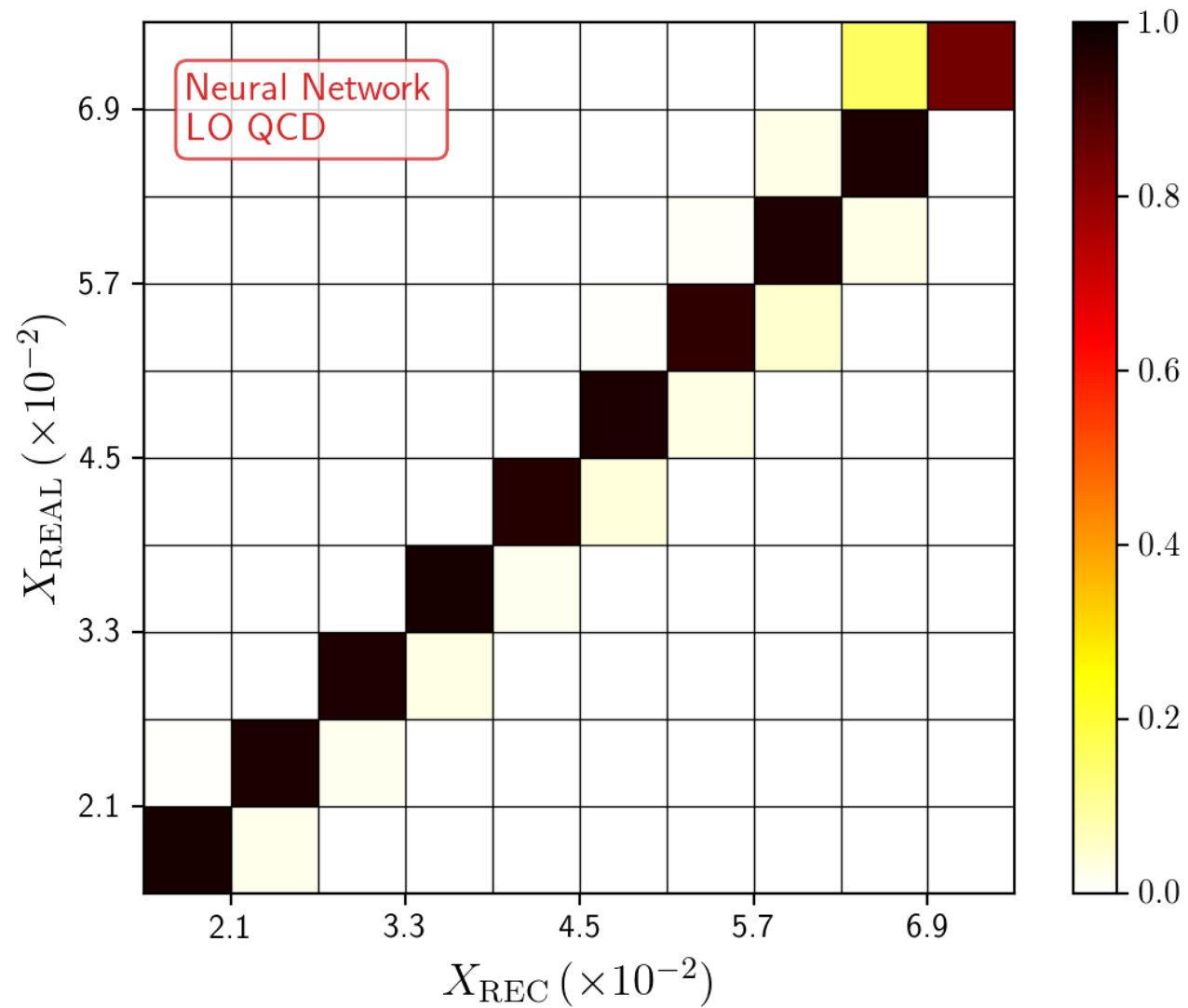
- Similar results for the reconstruction of z .





- ▶ Pass from one layer to the next by applying a non-linear **activation function** to a **weighted** sum of the previous layers.
- ▶ Pros: no need to play with the basis, less human bias (but not zero!)
- ▶ Cons: the complexity of the architecture requires more time for training.
- ▶ Also, one needs to choose the architecture.

Neural networks



- For LO the complexity of the NN greatly surpasses the complexity of the problem.

Going greener with ML

- ▶ All the running of MC codes takes a long time to reach good precision.
- ▶ Apart from boring, they carry a significant environmental impact (and to our pockets given the cost of electricity).
- ▶ What can we do to make things faster using current available resources?
 - ▶ Improve the codes: has to be done code by code.

Going greener with ML

- ▶ All the running of MC codes takes a long time to reach good precision.
- ▶ Apart from boring, they carry a significant environmental impact (and to our pockets given the cost of electricity).
- ▶ What can we do to make things faster using current available resources?
 - ▶ Improve the codes: has to be done code by code.
- ▶ Can we speed up the running time without touching the code?

Going greener with ML

- ▶ All the running of MC codes takes a long time to reach good precision.
- ▶ Apart from boring, they carry a significant environmental impact (and to our pockets given the cost of electricity).
- ▶ What can we do to make things faster using current available resources?
 - ▶ Improve the codes: has to be done code by code.
- ▶ Can we speed up the running time without touching the code?

Perhaps.

- ▶ Most codes require non perturbative inputs (e.g. PDFs).
- ▶ PDFs are provided as grids and functions that read the grids and interpolate over them (e.g. LHAPDF).
- ▶ And this is quite efficient, as long as we don't need to run millions and millions of calculations.

- ▶ Most codes require non perturbative inputs (e.g. PDFs).
- ▶ PDFs are provided as grids and functions that read the grids and interpolate over them (e.g. LHAPDF).
- ▶ And this is quite efficient, as long as we don't need to run millions and millions of calculations.
- ▶ A quick exploration shows that the time spent on the interpolation could be reduced 40-50% if we had analytical expressions for the PDFs.

- ▶ Most codes require non perturbative inputs (e.g. PDFs).
- ▶ PDFs are provided as grids and functions that read the grids and interpolate over them (e.g. LHAPDF).
- ▶ And this is quite efficient, as long as we don't need to run millions and millions of calculations.
- ▶ A quick exploration shows that the time spent on the interpolation could be reduced 40-50% if we had analytical expressions for the PDFs.

Goal: find an analytical x and Q^2 form for a set of proton PDFs.

- ▶ We are working (for now) with HERAPDF2.0

Idea and first results

- ▶ For most PDFs the x dependence at some initial scale is written as

$$f_i(x, Q_0^2) = N_i x^{\alpha_i} (1 - x)^{\beta_i} P(x, c_{ij})$$

- ▶ We propose

$$f_i(x, Q_0^2, Q^2) = (N_i + g_{i,1}(Q^2, Q_0^2)) x^{\alpha_i + g_{i,2}(Q^2, Q_0^2)} (1 - x)^{\beta_i + g_{i,3}(Q^2, Q_0^2)} P(x, c_{ij} + g_{i,4}(Q^2, Q_0^2))$$

with $g_{i,j}(Q_0^2, Q_0^2) = 0$

In particular, for now, we're exploring $g_{i,4}(Q^2, Q_0^2) = 0$

- ▶ With that simplification, the ratio of the same flavour PDF at different scales is

$$R_i(x, Q_0^2, Q^2) \propto x^{g_{i,2}(Q^2, Q_0^2)}(1 - x)^{g_{i,3}(Q^2, Q_0^2)}$$

- ▶ Taking logarithm

$$\ln(R_i) = \ln\left(\frac{N'_i(Q^2, Q_0^2)}{N_i}\right) + g_{i,2}(Q^2, Q_0^2)\ln(x) + g_{i,3}(Q^2, Q_0^2)\ln(1 - x)$$

- ▶ With that simplification, the ratio of the same flavour PDF at different scales is

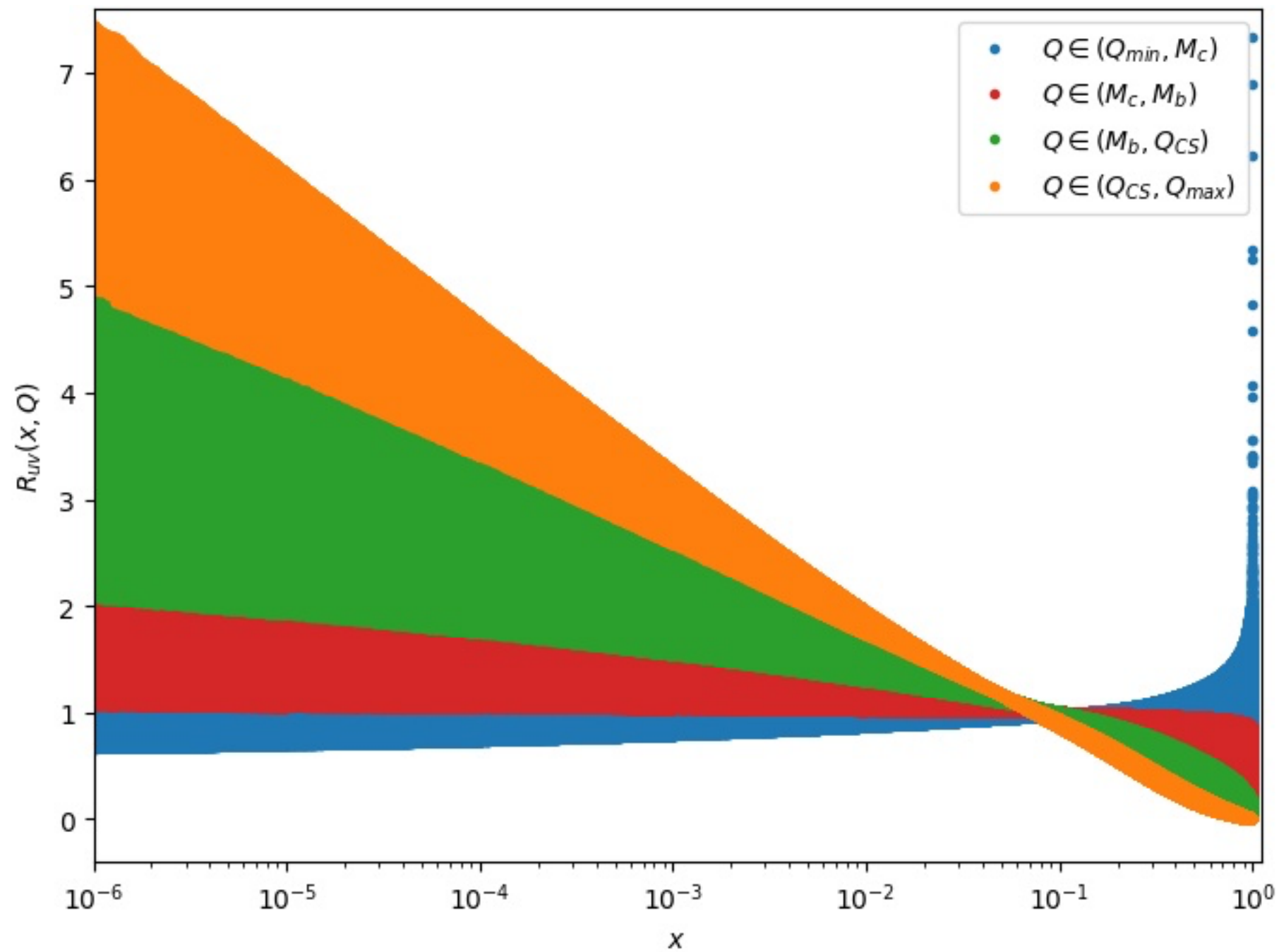
$$R_i(x, Q_0^2, Q^2) \propto x^{g_{i,2}(Q^2, Q_0^2)}(1 - x)^{g_{i,3}(Q^2, Q_0^2)}$$

- ▶ Taking logarithm

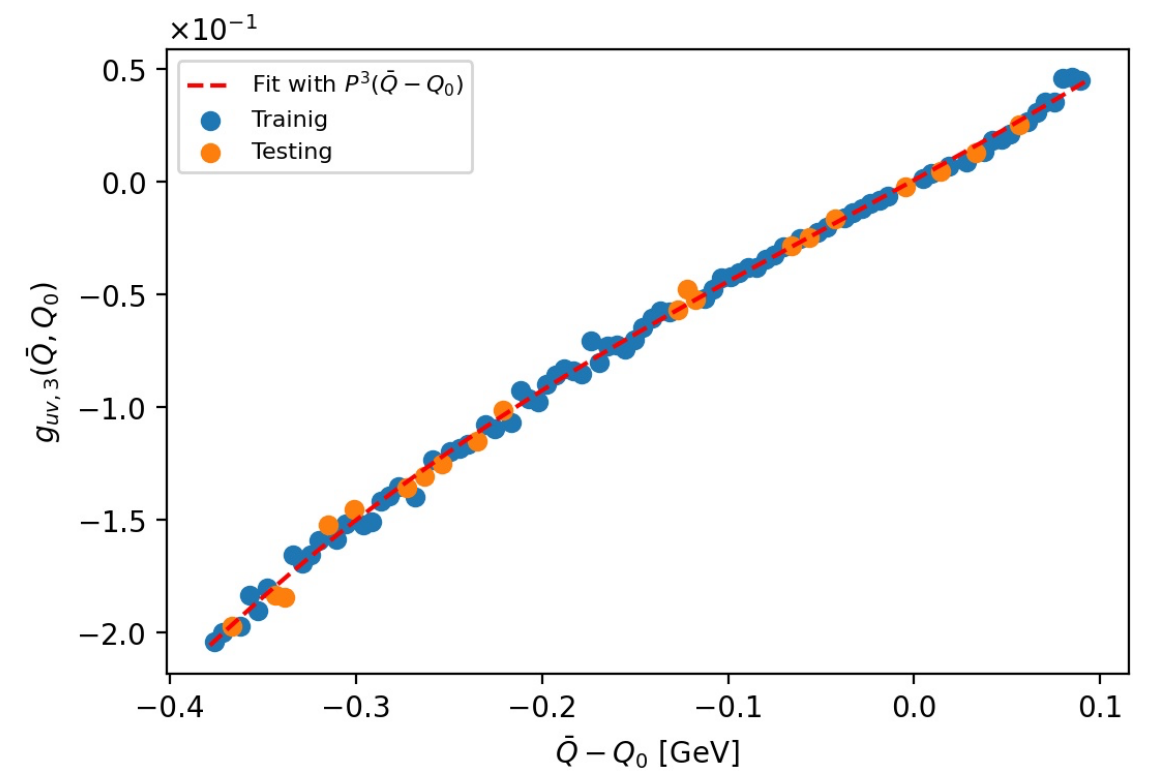
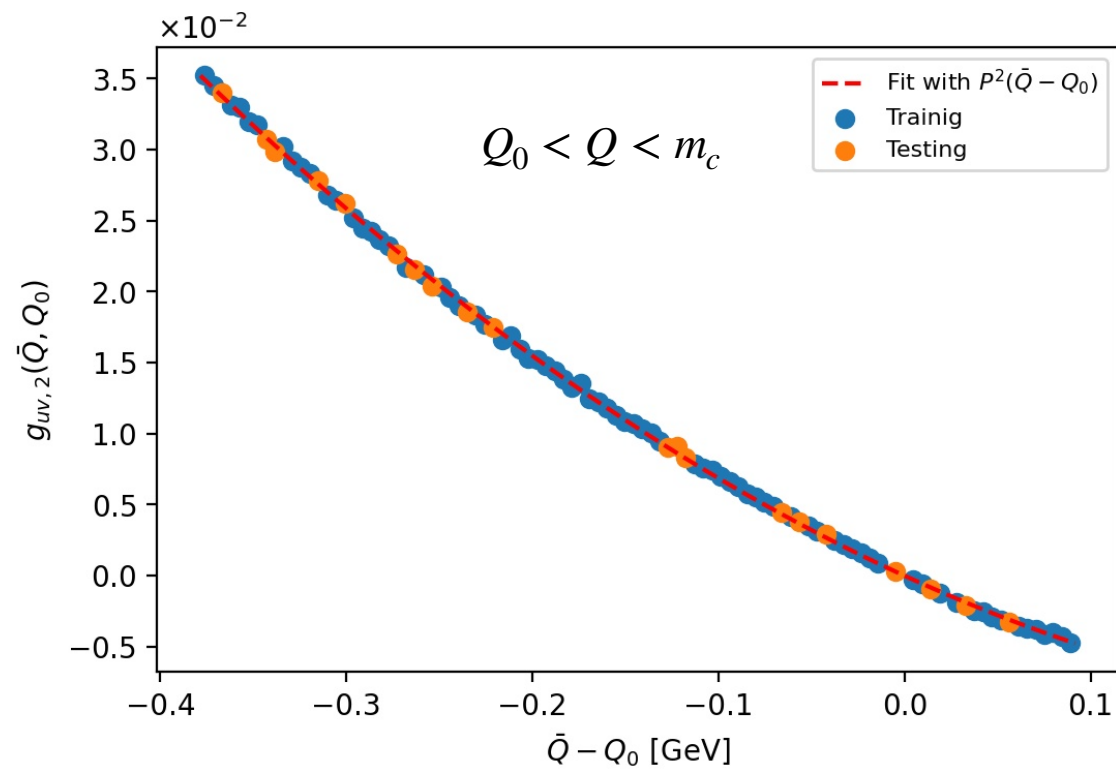
$$\ln(R_i) = \ln\left(\frac{N'_i(Q^2, Q_0^2)}{N_i}\right) + g_{i,2}(Q^2, Q_0^2)\ln(x) + g_{i,3}(Q^2, Q_0^2)\ln(1 - x)$$

- ▶ Now we only have to find the missing functions.

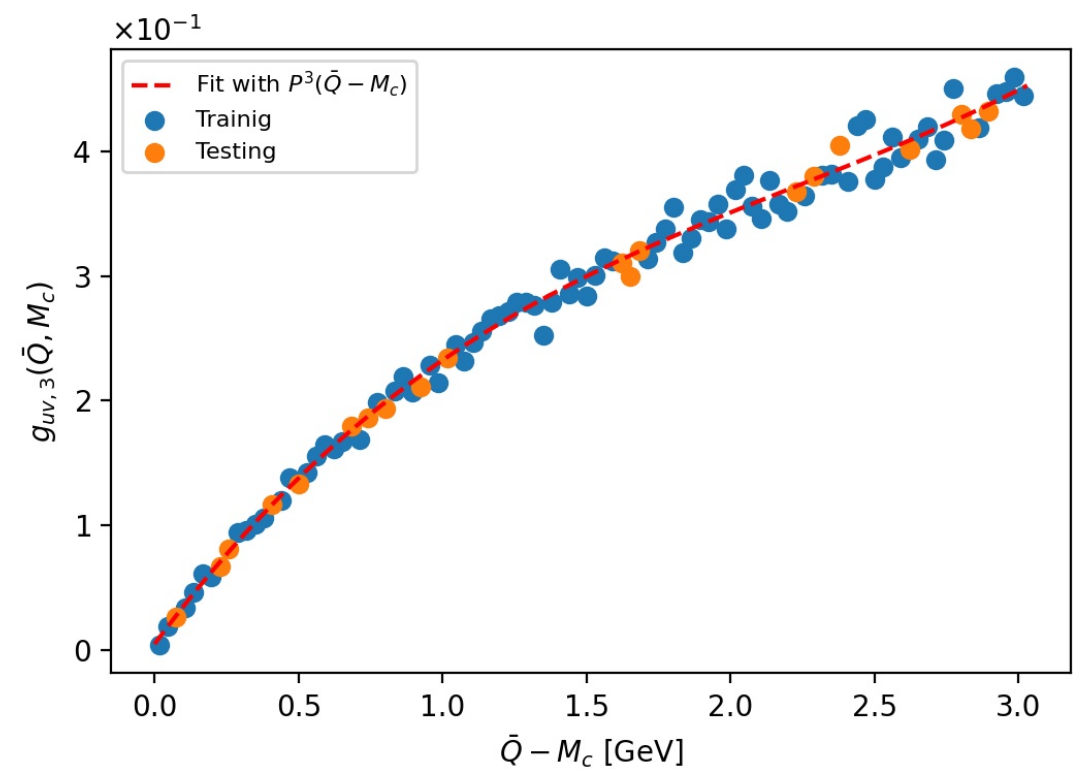
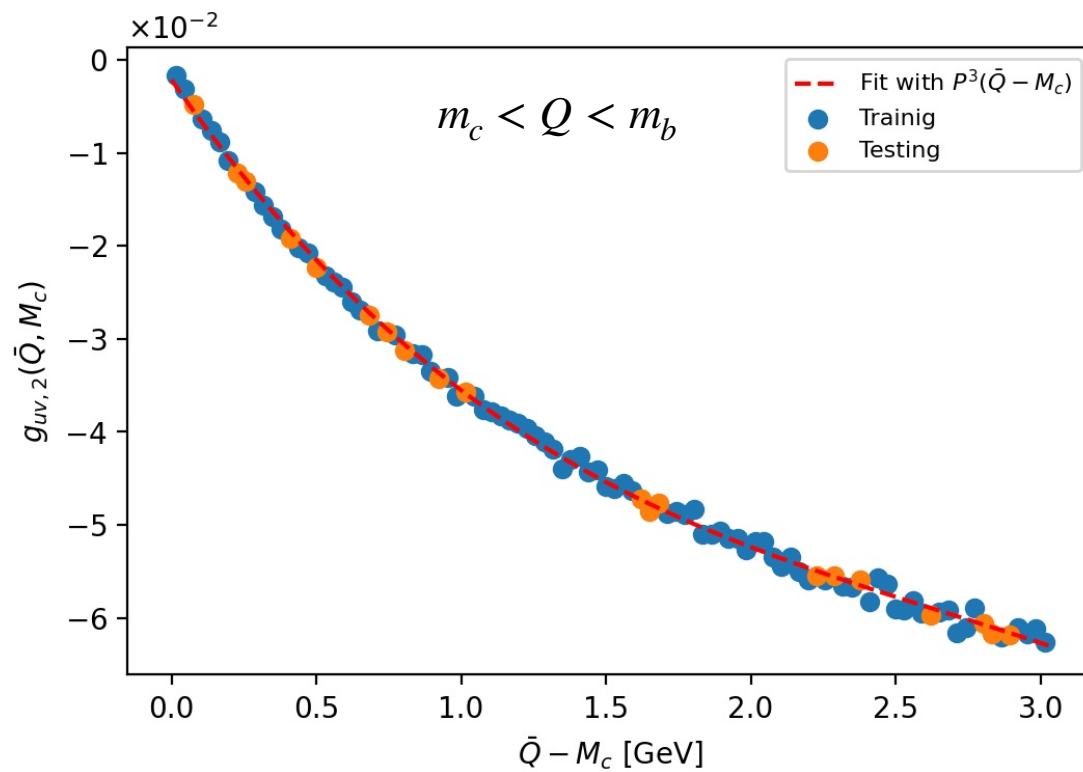
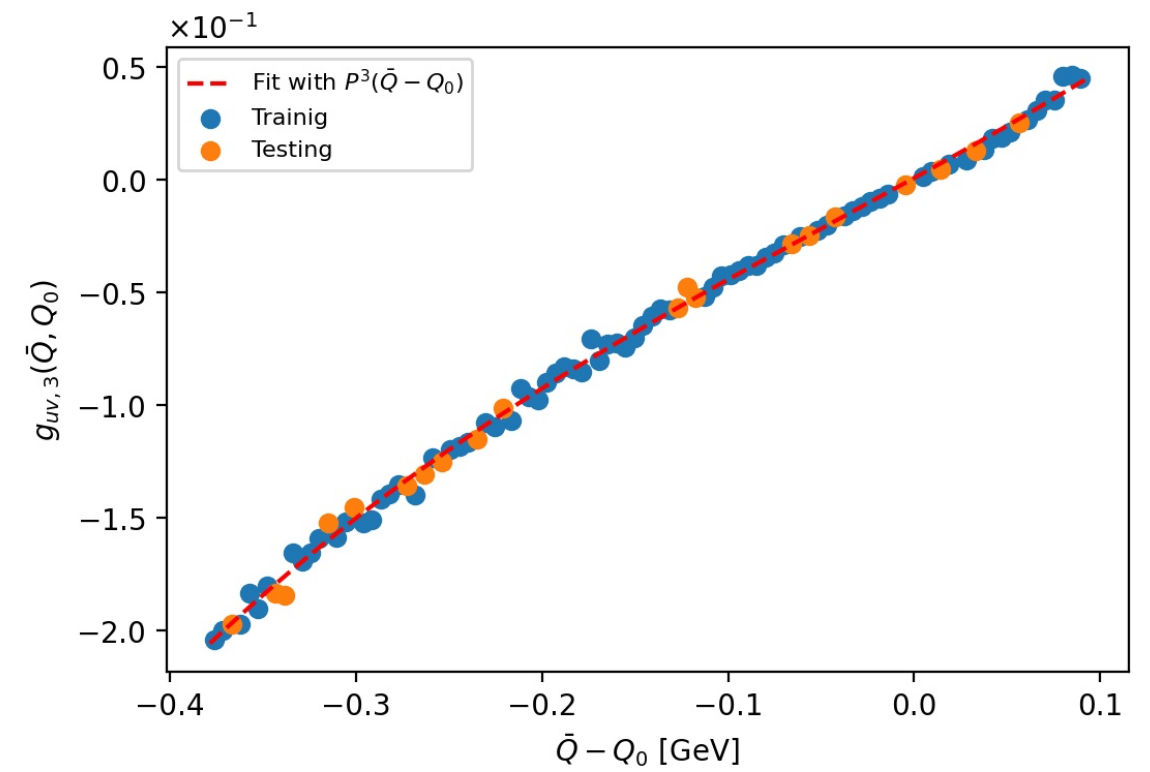
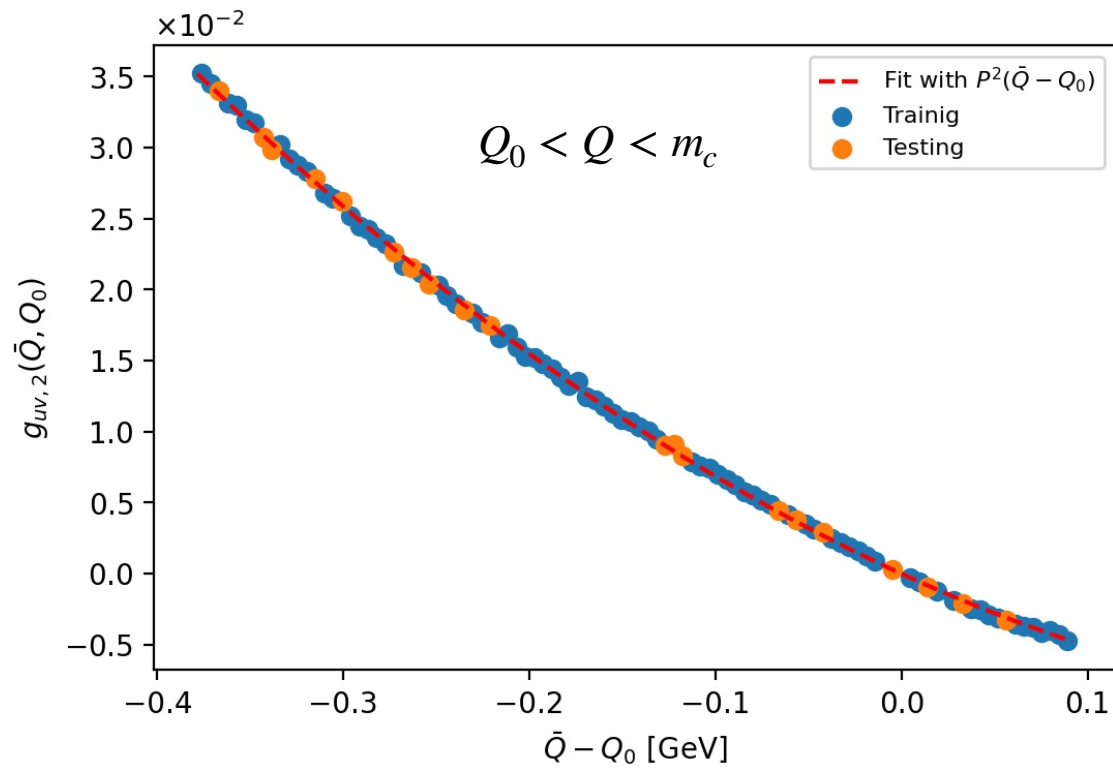
First results: valence up



First results: valence up



First results: valence up



The GAPP Initiative

- ▶ This work is part of the GAPP Initiative.
- ▶ We aim to quantify the carbon footprint of HEP research and study ways of reducing it.
- ▶ If you are doing phenomenological studies and would like to contribute to the project, please send an email to gapp-initiative@googlegroups.com with:
 - ▶ **brief description of the simulation**
 - ▶ **time needed to run**
 - ▶ **hardware used**

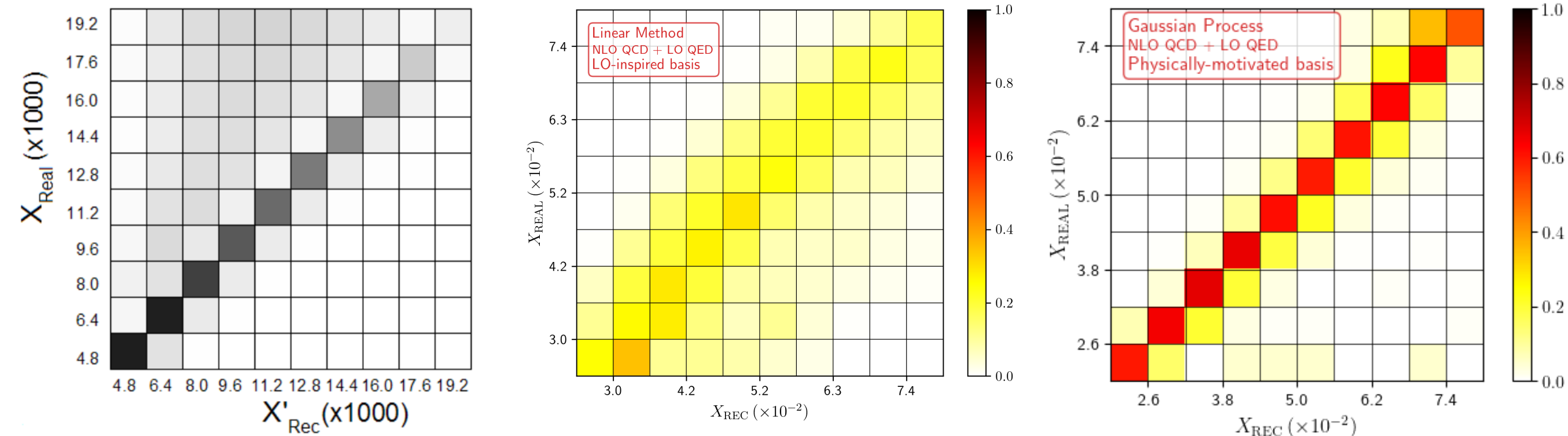
The GAPP Initiative

- ▶ This work is part of the GAPP Initiative.
- ▶ We aim to quantify the carbon footprint of HEP research and study ways of reducing it.
- ▶ If you are doing phenomenological studies and would like to contribute to the project, please send an email to gapp-initiative@googlegroups.com with:
 - ▶ **brief description of the simulation**
 - ▶ **time needed to run**
 - ▶ **hardware used**

Join the GAPP!

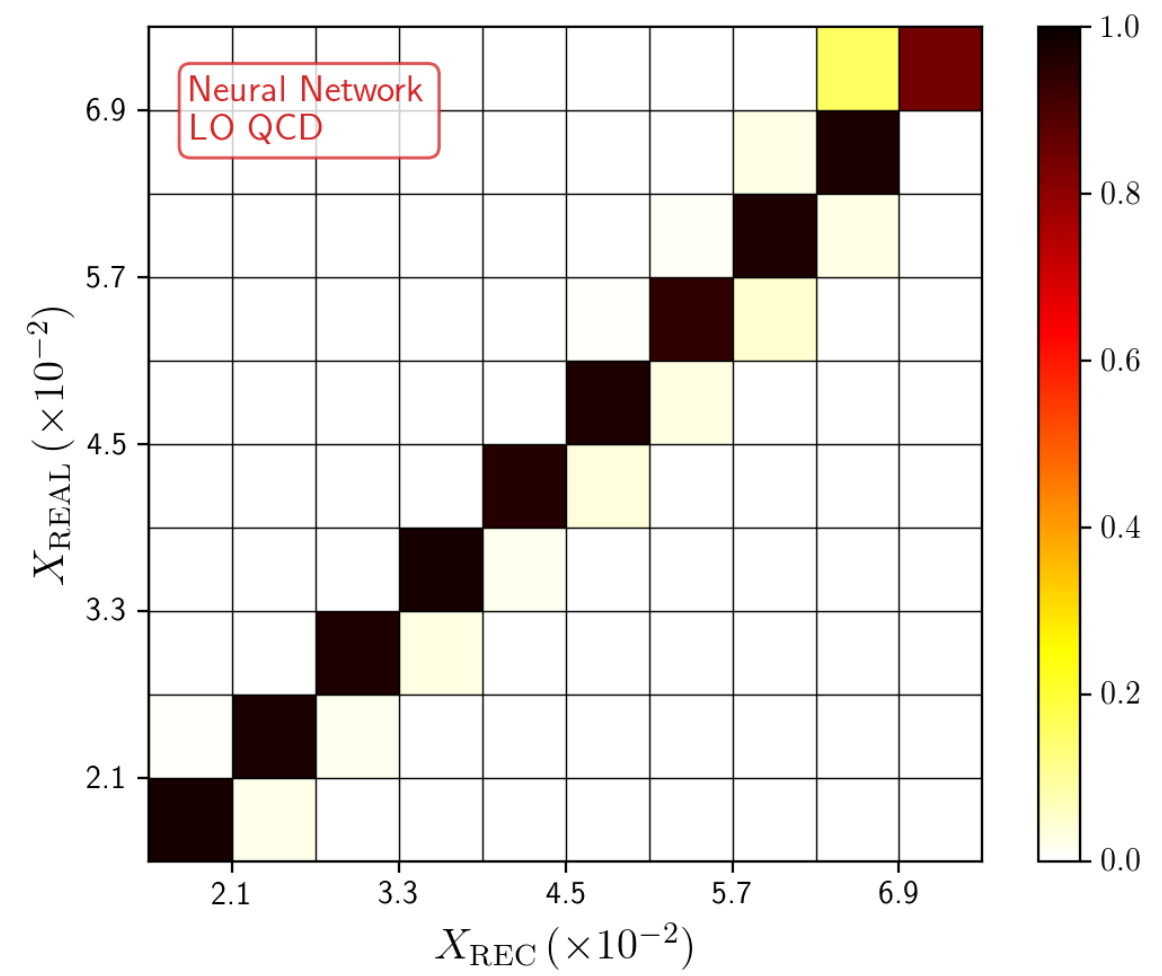
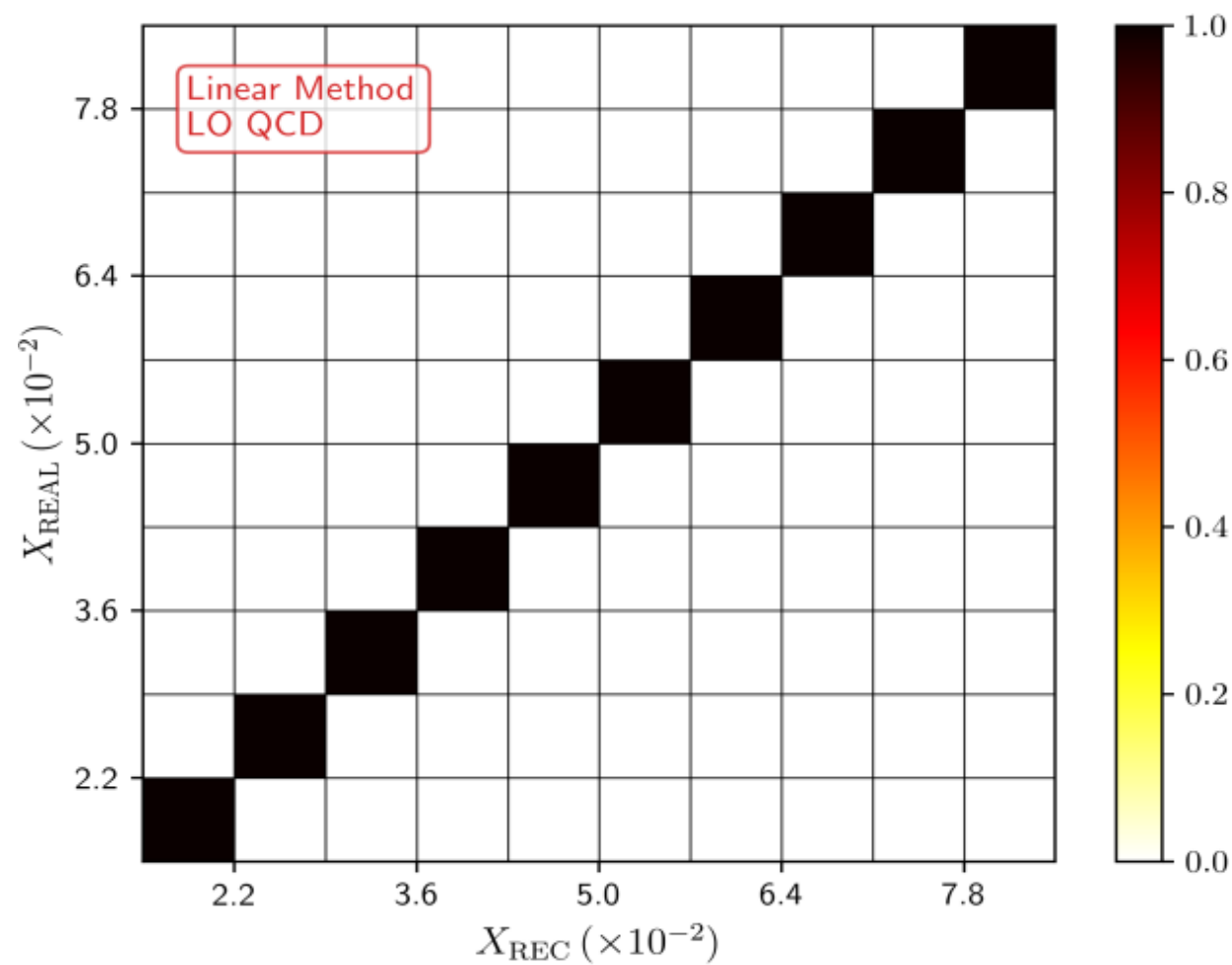
Summary

- ▶ We have explored the applicability of ML techniques to better understand the underlying kinematics of a p+p collision.
- ▶ The methodology can be used for any process by non-experts.
- ▶ The methods applied can result in a better reconstruction than in the original work, but physical intuition can play a relevant role.

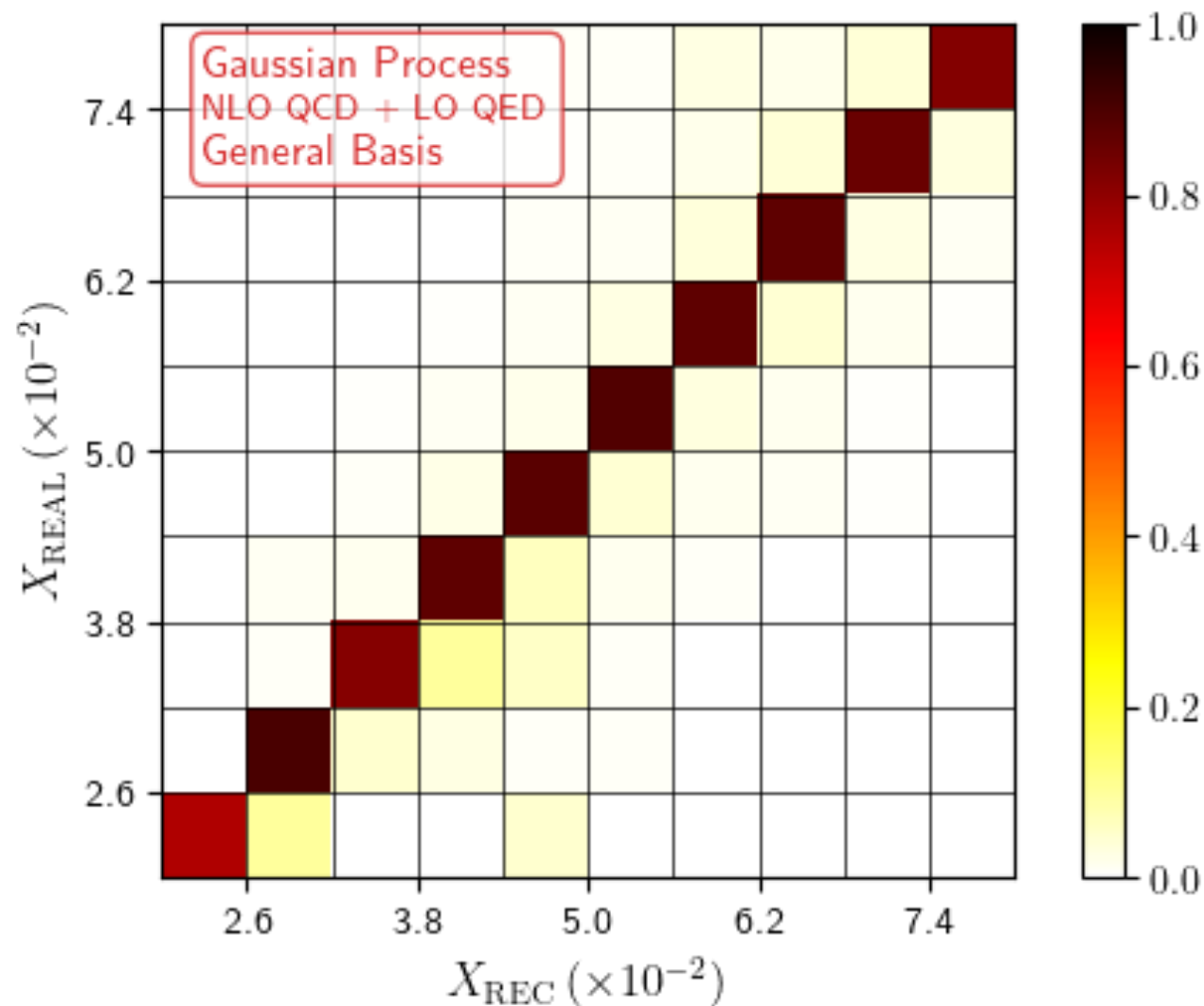


- Higher sophistication of the method does not always translate into better results. E.g.:

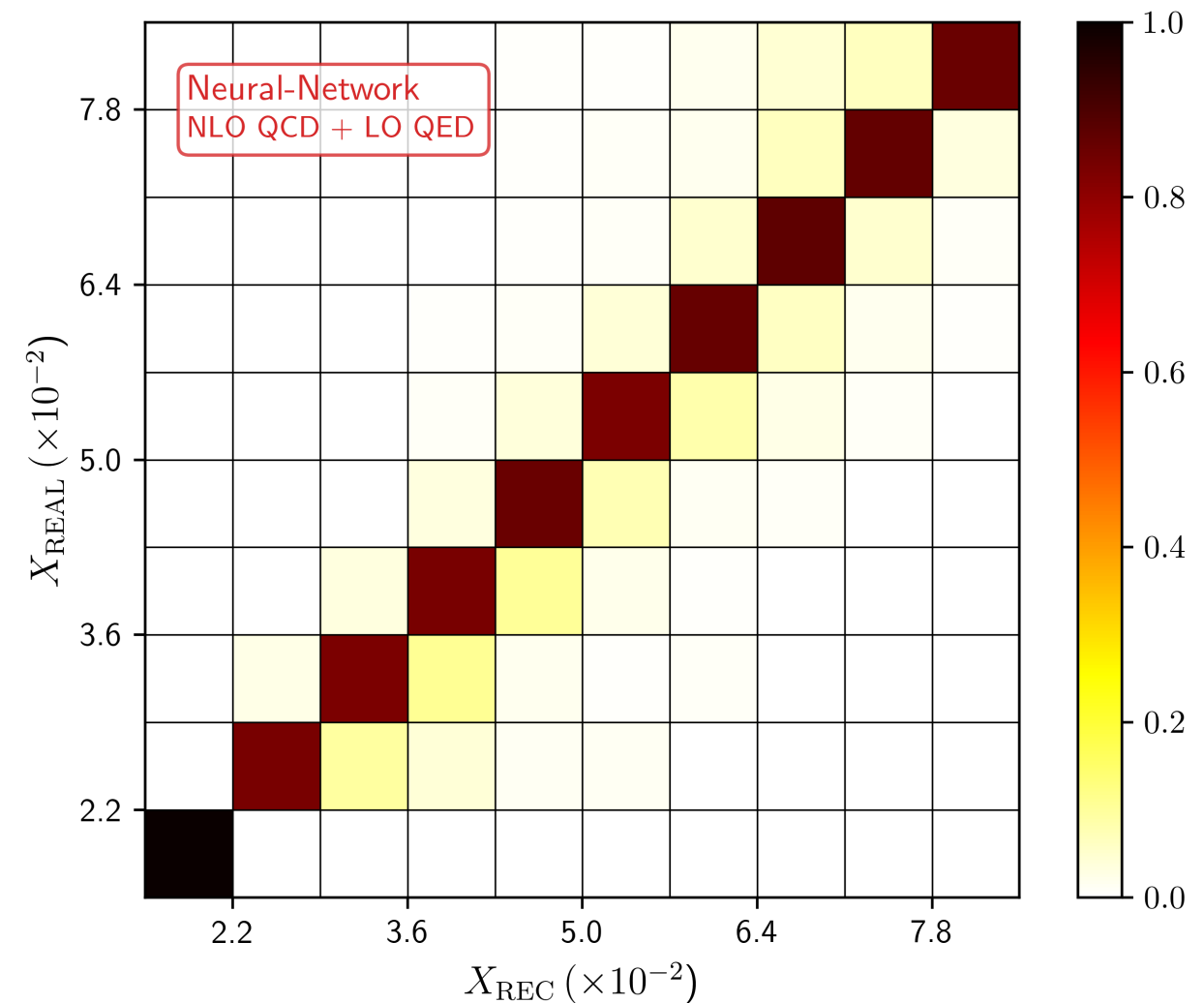
exact relation known



- Higher sophistication of the method does not always translate into better results. E.g.:

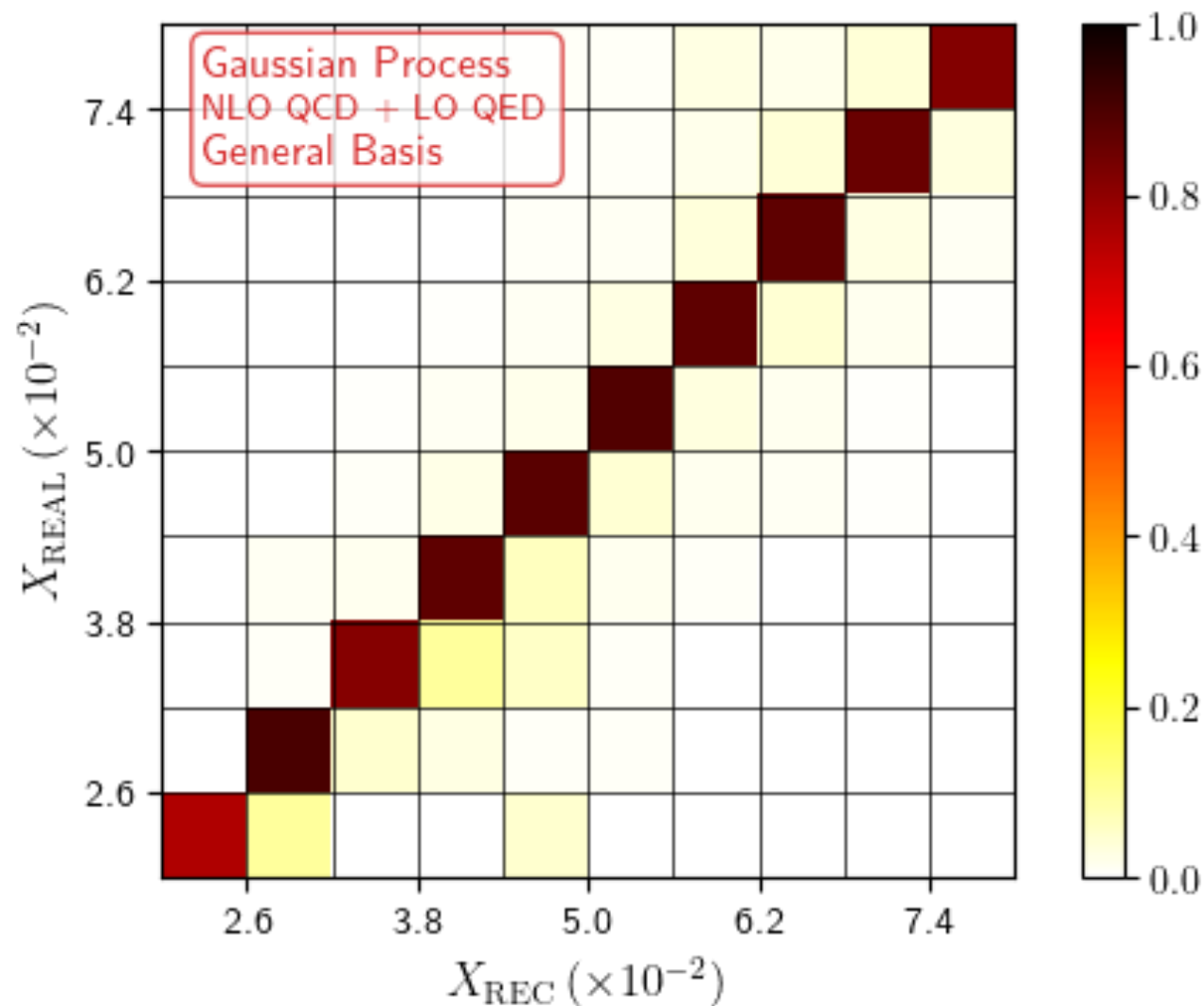


few minutes to train

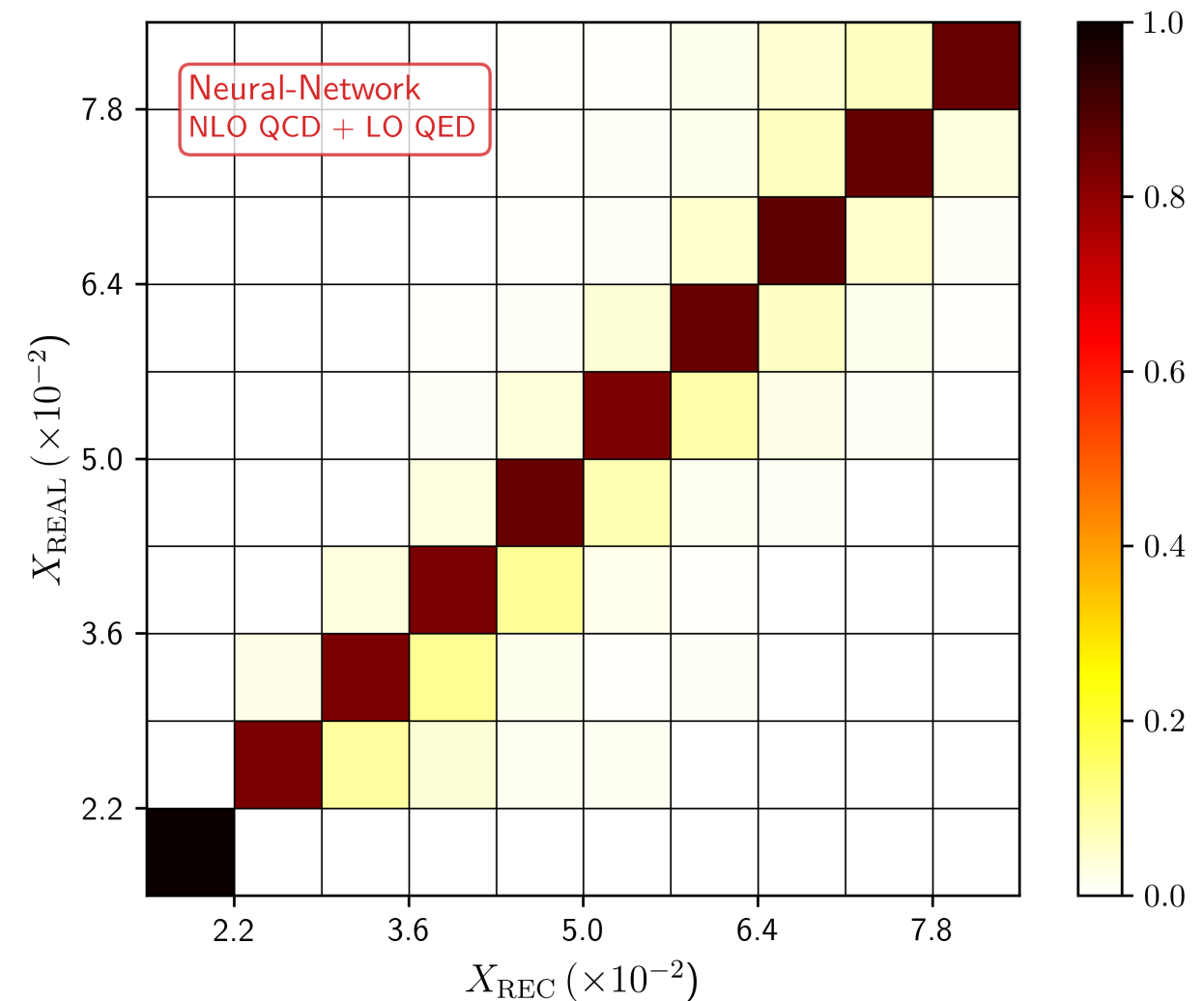


several hours to train
exploration of "good" architecture

- Higher sophistication of the method does not always translate into better results. E.g.:



few minutes to train



several hours to train
exploration of "good" architecture

- Promising steps in speeding up the calculation of codes using PDFs.

**Thank you for
your attention!**